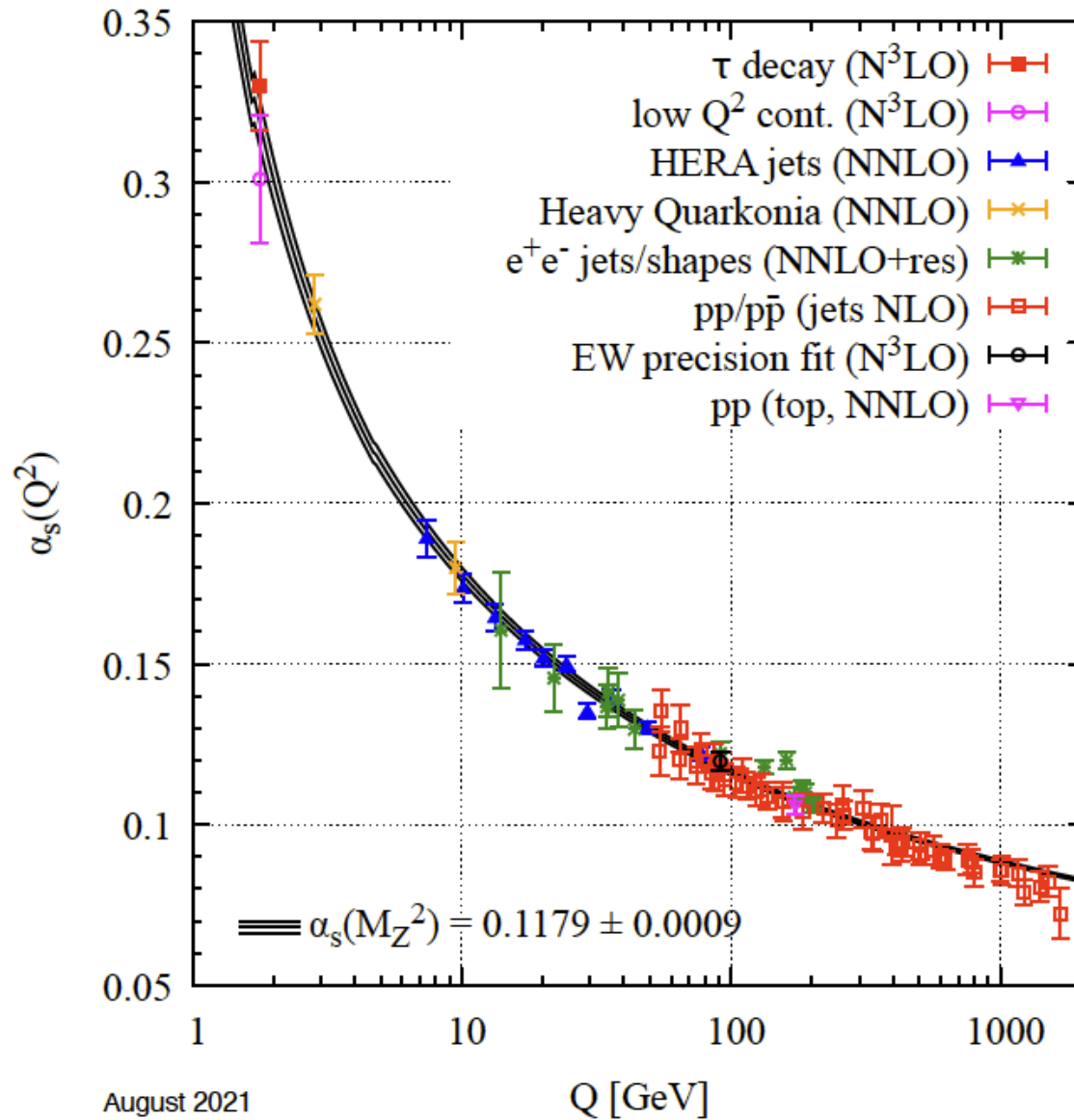


α_s measurements

- α_s measurements
- Deep Inelastic Scattering:
 - ★ Sum rules
 - ★ **Scaling violation**
 - ★ Jets in DIS
- **τ lepton decay**
- Quarkonium
- e^+e^- annihilation
 - ★ hadronic cross sections
 - ★ Z decay
 - ★ Event shapes
 - ★ Scaling violation in fragmentation
- **Hadron-Hadron scattering**

α_s measurements



August 2021

Figure 9.4: Summary of measurements of α_s as a function of the energy scale Q . The respective degree of QCD perturbation theory used in the extraction of α_s is indicated in brackets (NLO: next-to-leading order; NNLO: next-to-next-to-leading order; NNLO+res.: NNLO matched to a resummed calculation; N³LO: next-to-NNLO).

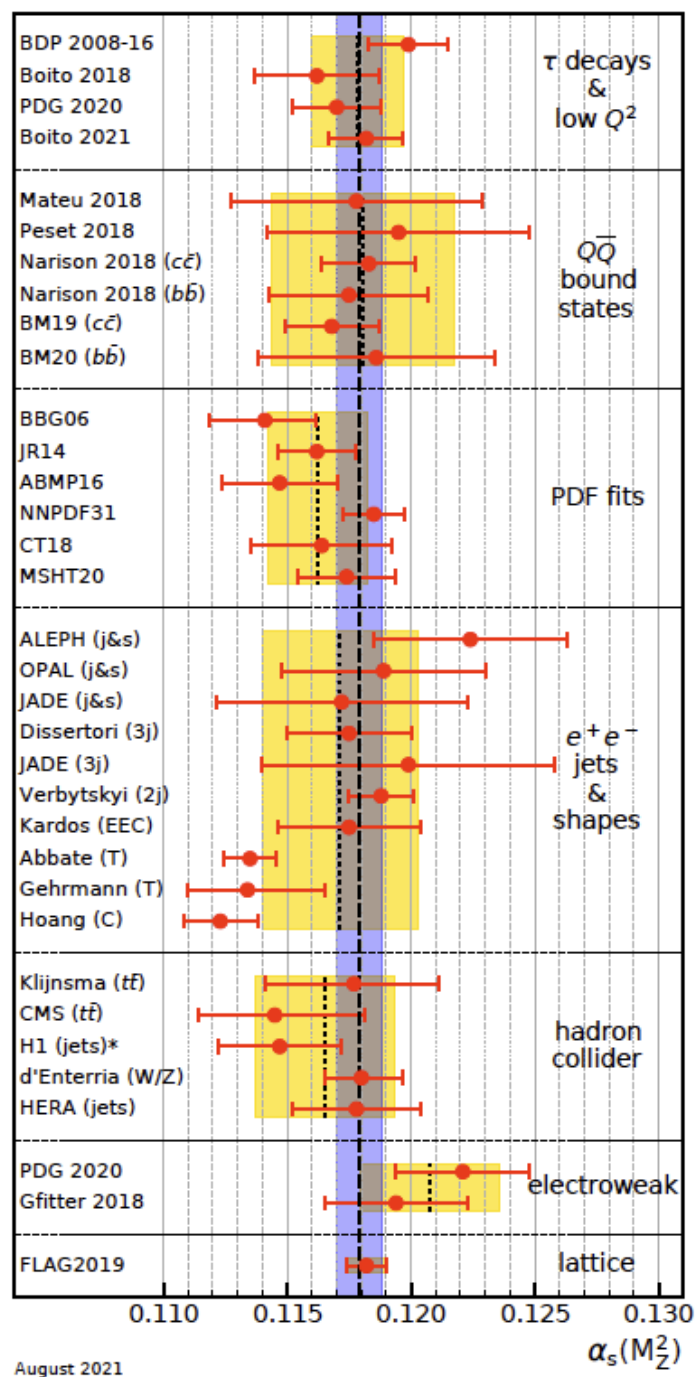


Figure 9.2: Summary of determinations of $\alpha_s(M_Z^2)$ from the seven sub-fields discussed in the text. The yellow (light shaded) bands and dotted lines indicate the pre-average values of each sub-field. The dashed line and blue (dark shaded) band represent the final world average value of $\alpha_s(M_Z^2)$. The “*” symbol within the “hadron colliders” sub-field indicates a determination including a simultaneous fit of PDFs.

α_s measurement: Scaling Violation in DIS

- ★ In the evolution of F_2 at intermediate value of x ($0.01 < x < 0.3$), $\alpha_s(M_Z)$ and the $xg(x)$ are strongly correlated through the DGLAP. An increase of $\alpha_s(M_Z)$ can be compensated by a harder $xg(x)$. This has limited, in the past, the precision of the $\alpha_s(M_Z)$ measurements from NLO DGLAP fits using DIS data.
- ★ At small x , ($x < 0.01$) this correlation becomes weaker $\Rightarrow xg(x)$ drives the behaviour of F_2 (not only of $dF_2/d\ln Q^2$).
- ★ The jets data help a lot: **BGF** directly sensible to $xg(x)$, **QCDC** to $\alpha_s(M_Z)$
- ★ A simultaneous fit of $\alpha_s(M_Z)$ and of the parameters of the p.d.f. gives:

$$\alpha_s(M_Z) = 0.1156 \pm 0.0011 \text{ (exp.)}_{-0.0002}^{+0.0001} \text{ (model+parameterization)} \pm 0.0029 \text{ (scale)}$$

to compare with $\alpha_s(M_Z) = 0.1179 \pm 0.0009$ (PDG).

α_s measurement: τ lepton decay

- $R_\tau^{th} = \frac{\Gamma(\tau \rightarrow \nu_\tau + \text{hadrons})}{\Gamma(\tau \rightarrow \nu_\tau + \bar{\nu}_e + e)} = 3S_{EW}(1 + \delta_{pQCD} + \delta_{NPQCD})$

- $S_{EW} = 1.0194$ from the electroweak theory

- $\delta_{pQCD} = \frac{\alpha_s(m_\tau^2)}{\pi} + 5.2 \left(\frac{\alpha_s(m_\tau^2)}{\pi} \right)^2 + 26.4 \left(\frac{\alpha_s(m_\tau^2)}{\pi} \right)^3 \left(\pm 130 \left(\frac{\alpha_s(m_\tau^2)}{\pi} \right)^4 \right)$

- $\delta_{NPQCD} = -0.007 \pm 0.004$

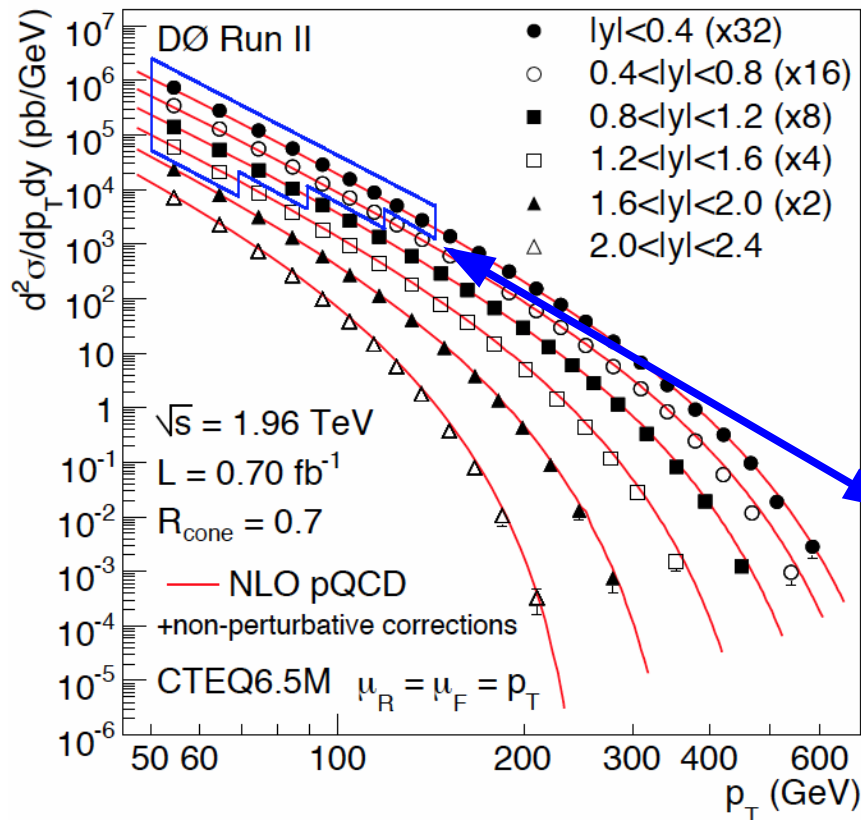
- $R_\tau^{exp} = \frac{1 - B_e - B_\mu}{B_l} = \frac{1}{B_l} - 1 - f_\mu$ with $f_\mu = 0.9726$ phase space correction

B_l can be measured in 3 different manners: from the leptonic branching ratio of e, μ , $B_e = B_l e B_\mu = f_\mu B_l$, and from the life time $\tau_\tau = B_l \tau_\mu (m_\mu/m_\tau)^5$

- $\alpha_s(M_\tau^2) = 0.312 \pm 0.015 \rightarrow \alpha_s(M_Z^2) = 0.1178 \pm 0.0019$

α_s measurement: hadron collider jets

- D0, the other experiment at the TEVATRON collider, has measured α_s with the Run II data
- The coupling is obtained by the single-jet inclusive cross section at a CM energy of $\sqrt{s}=1.96$ TeV
- The $d^2\sigma_{\text{jet}}/dp_T dy$ is measured (y is the rapidity)
- The jets are recognized by a cone-algorithm with radius $R = 0.7$



D0 uses only these events with central jets, where the detector is well calibrated $|y| < 1.6$, and at an energy not very high ($50 < p_T < 145$ GeV) to avoid a region where the cross section for jet production is used to constrain the PDF.

α_s measurement: hadron collider jets

- From the theoretical point of view:

$$\sigma(\alpha_s(M_Z)) = \sigma_{\text{pert}}(\alpha_s(M_Z)) \cdot C_{\text{nonpert}} = \left(\left(\sum_n \alpha_s^n c_n \right) \otimes f_1(\alpha_s) \otimes f_2(\alpha_s) \right) \cdot C_{\text{nonpert}}$$

- f_1, f_2 are the structure functions of the proton and of the antiproton
- \otimes convolution,
- C_{nonpert} contains the corrections due to the hadronization and underlying event
- The α_s value is obtained through a minimization procedure between experimental data and the theoretical formula.

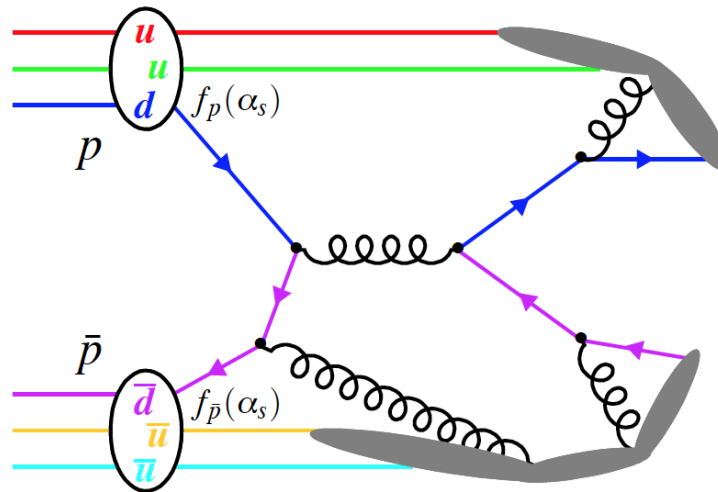
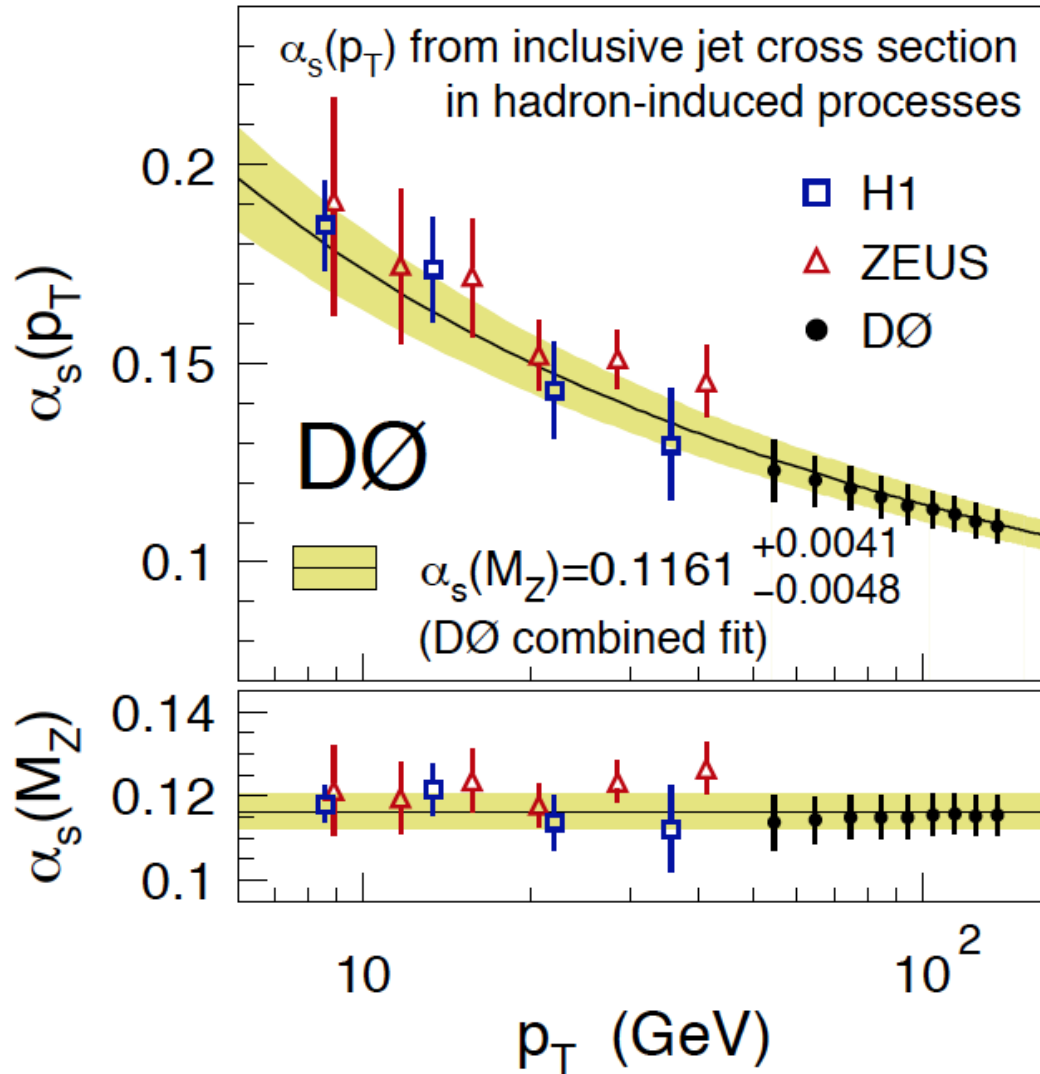


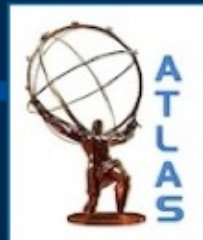
Figure 1: A drawing, illustrating jet production in a proton anti-proton collision with the hard scattering process, initial state, final state radiation and hadronisation (jet fragmentation) including the underlying event.

α_s measurement: hadron collider jets



$$\alpha_s(M_Z) = 0.1161^{+0.0041}_{-0.0048}$$

Inclusive Jet Ratios



Inclusive 3-Jet / Inclusive 2-Jet Ratio ($R_{3/2}, N_{3/2}$) ($\sqrt{s} = 7 \text{ TeV}$)

$$R_{3/2} = \frac{\sigma_{3\text{-jet}}}{\sigma_{2\text{-jet}}}$$

(lead jet)

$$N_{3/2} = \frac{\sigma_{3\text{-jet}}}{\sigma_{2\text{-jet}}}$$

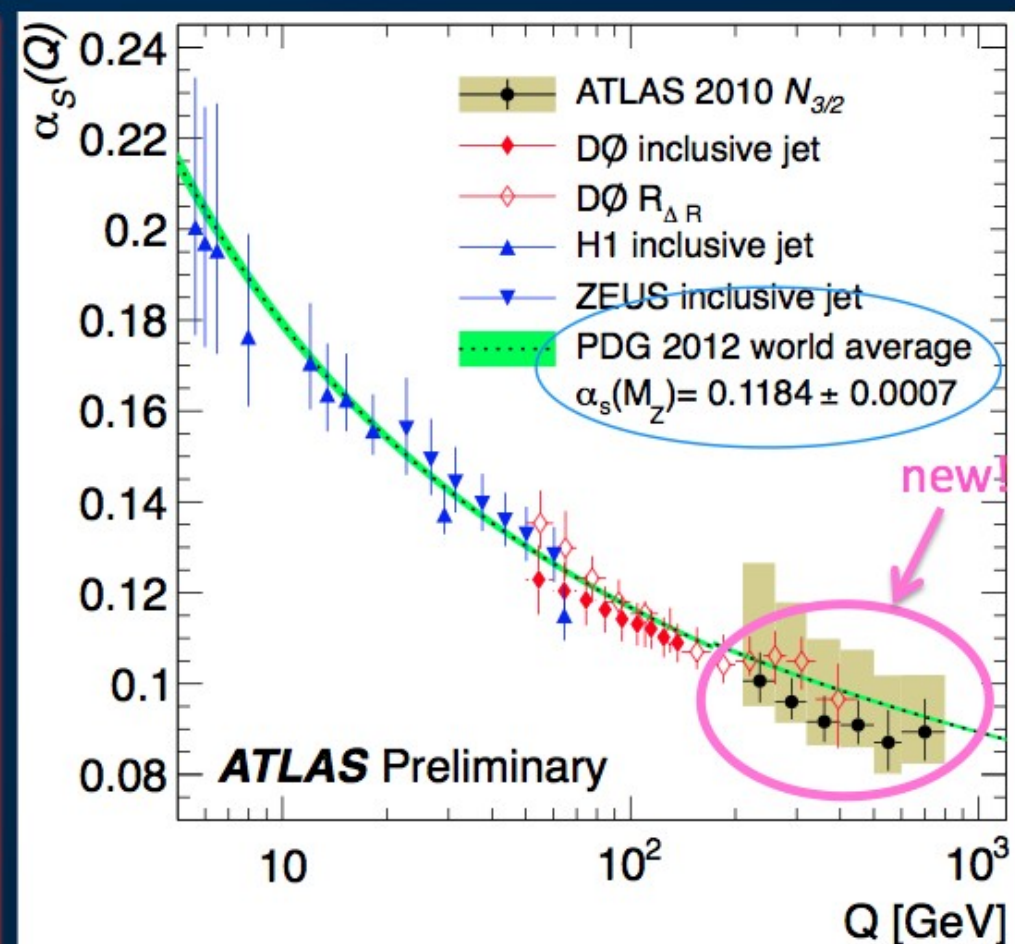
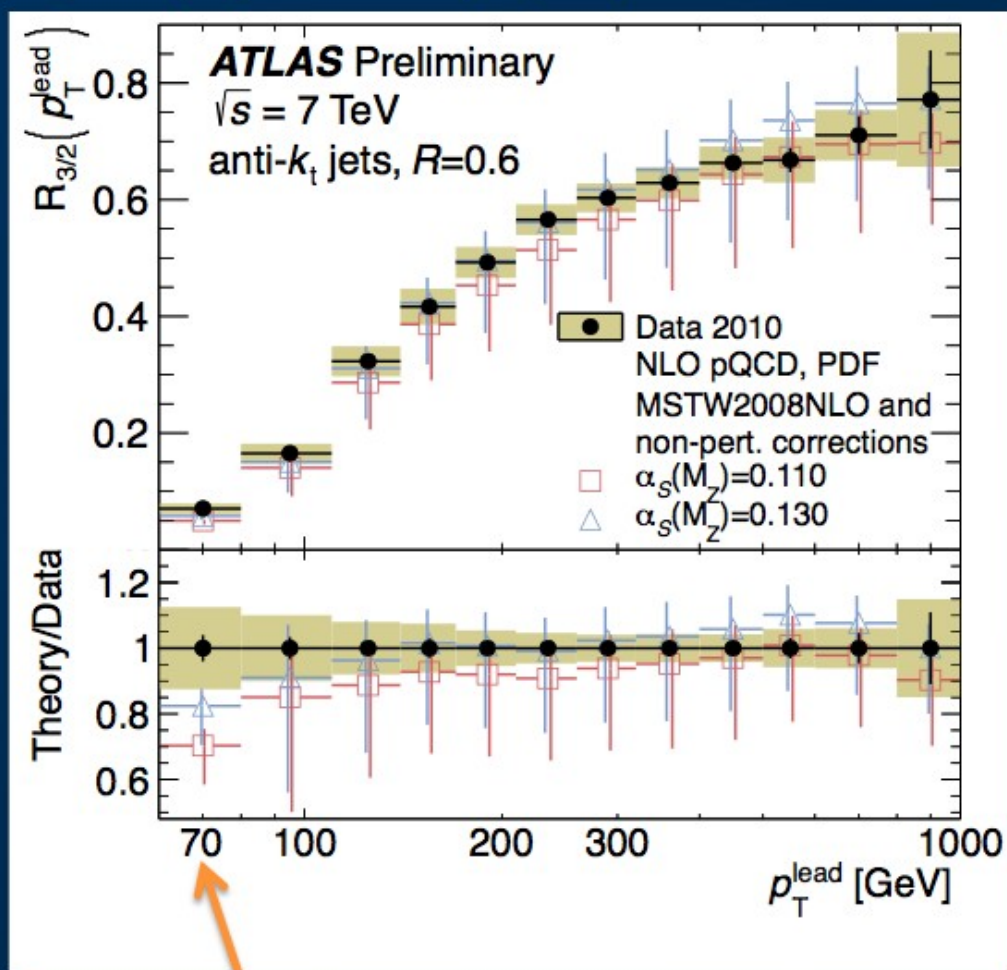
(all jets)



α_s

▶ Extract α_s :

$$\alpha_s(M_Z) = 0.111 \pm 0.006 \text{ (exp.) } {}^{+0.016}_{-0.003} \text{ (theory)}$$



▶ Good agreement overall, except here for $p_T^{\text{lead}} < 140 \text{ GeV}$

ATLAS: ATLAS-CONF-2013-041

Inclusive Jet Ratios



Inclusive 3-Jet / Inclusive 2-Jet Ratio ($R_{3/2}$) ($\sqrt{s} = 7$ TeV)

Measure $R_{3/2}$ vs. $\langle p_{T1,2} \rangle$, the average p_T of the two leading jets in the event

Compare data to NLO QCD \rightarrow Good agreement!

Extract

$$\alpha_s(M_Z) = 0.1148 \pm 0.0014 \text{ (exp.)} \pm 0.0018 \text{ (PDF)} {}^{+0.0050}_{-0.0000} \text{ (scale)}$$

$$250 < \langle p_{T1,2} \rangle < 1390 \text{ GeV}$$

► First determination of α_s from measurements at Q scales up to 1 TeV

