



Università degli Studi di Padova

LESSON 17: REACTION OF SOLUTE WITH DISPERSION





In estuaries and along coastal plain rivers the transport of the solute/pollutant strongly depends also on dispersion. The equation of the transport is the following:

$$\frac{\partial C}{\partial t} + \frac{Q}{A}\frac{\partial C}{\partial x} = \frac{1}{A}\frac{\partial}{\partial x}\left(K_x\frac{\partial C}{\partial x}\right) \mp rC$$

The dynamics of \tilde{O} and BOD is then described by the following system of equations:

$$\begin{pmatrix} \frac{\partial B}{\partial t} + \frac{Q}{A} \frac{\partial B}{\partial x} = \frac{1}{A} \frac{\partial}{\partial x} \left(K_x A \frac{\partial B}{\partial x} \right) - r_d B \\ \frac{\partial \tilde{O}}{\partial t} + \frac{Q}{A} \frac{\partial \tilde{O}}{\partial x} = \frac{1}{A} \frac{\partial}{\partial x} \left(K_x A \frac{\partial \tilde{O}}{\partial x} \right) + r_d B - r_a \tilde{O} \end{cases}$$

The solution of the full equations is only numerical. However an analytical solution is possible under the following assumptions:

- Stationarity, i.e. $\partial/\partial t = 0$
- Homogeneous reach, i.e. K_x , Q and A are constant

$$\longrightarrow U = \frac{Q}{A}$$

The two equations of the system now read:









$$\begin{cases} \frac{\mathrm{d}^2 B}{\mathrm{d}x^2} - \frac{U}{K_x} \frac{\mathrm{d}B}{\mathrm{d}x} - \frac{r_d}{K_x} B = 0 \qquad (1)\\ \frac{\mathrm{d}^2 \tilde{O}}{\mathrm{d}x^2} - \frac{U}{K_x} \frac{\mathrm{d}\tilde{O}}{\mathrm{d}x} - \frac{r_a}{K_x} \tilde{O} + \frac{r_d}{K_x} B = 0 \qquad (2) \end{cases}$$

Let's solve eq. (1). It is an homogenous II order differential equation, that in the most general form reads:

$$\frac{\mathrm{d}^2 y}{\mathrm{d}x^2} - a_1 \frac{\mathrm{d}y}{\mathrm{d}x} - a_0 y = 0$$

Its solution is given by the characteristic polynomial: $P(\lambda) = \lambda^2 - a_1\lambda - a_0 = 0$ The solution is then:

$$\lambda_{1,2} = \frac{a_1 \pm \sqrt{a_1^2 + 4a_0}}{2} \qquad \longrightarrow \qquad \lambda_{1,2} = \frac{1}{2} \left(\frac{U}{K_x} \pm \sqrt{\frac{U^2}{K_x^2} + \frac{4r_d}{K_x}} \right)$$

By rearranging the roots solution:







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BOD DISTRIBUTION



$$\lambda_{1,2} = \frac{1}{2} \left(\frac{U}{K_x} \pm \frac{U}{K_x} \sqrt{1 + \frac{4r_d K_x}{U^2}} \right) \longrightarrow \lambda_{1,2} = \frac{U}{2K_x} (1 \pm m_d)$$

ith: $m_d = \sqrt{1 + \frac{4r_d K_x}{U^2}} > 1$

Being $m_d > 0$ both λ_1 and λ_2 are real ($\lambda_1, \lambda_2 \in \mathcal{R}$), i.e. the solution of the equation is:

$$y = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \qquad \longrightarrow \qquad B = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x}$$

In this case $\lambda_1 > 0$ and $\lambda_2 < 0$, moreover $|\lambda_1| > |\lambda_2|$. The two constants of integration c_1 and c_2 can be defined as following.

Let's consider the discharge of *BOD* in x = 0. Because $K_x \neq 0$, *BOD* spreads also upstream (x < 0). The BCs are:

a.
$$x = 0 \longrightarrow B = B_0$$

b.
$$x \to +\infty \longrightarrow B = 0$$

$$C. \quad \chi \to -\infty \quad \longrightarrow \quad B = 0$$







BOD DISTRIBUTION

B

 B_0

x < 0

 $\lambda_1 > 0$

cusp



x > 0 $\lambda_2 < 0$

By applying the conditions (b) and (c), we find respectively:

$$B = c_1 e^{\lambda_1 \infty} + c_2 e^{\lambda_2 \infty} = 0 \quad \longrightarrow \quad c_1 = 0 \text{ for } x > 0$$
$$B = c_1 e^{-\lambda_1 \infty} + c_2 e^{-\lambda_2 \infty} = 0 \quad \longrightarrow \quad c_2 = 0 \text{ for } x < 0$$

We must split the solution:

 $\begin{cases} B = c_2 e^{\lambda_2 x} & x \ge 0 \\ B = c_1 e^{\lambda_1 x} & x < 0 \end{cases}$

Considering the condition (a) and expanding the equations:

$$\begin{cases} B = B_0 e^{\frac{U}{2K_x}(1-m_d)x} & x \ge 0\\ B = B_0 e^{\frac{U}{2K_x}(1+m_d)x} & x < 0 \end{cases}$$





Environmental Fluid Mechanics – Lesson 17: Reaction and Dispersion



X



EVALUATION OF B_0



 B_0 is calculated by the following mass balance:

$$QB_1(-\Delta x) - K_x A \frac{\mathrm{d}B_1}{\mathrm{d}x} \bigg|_{x = -\Delta x} + Q_w B_w(0) = QB_2(\Delta x) - K_x A \frac{\mathrm{d}B_2}{\mathrm{d}x} \bigg|_{x = \Delta x}$$

If $\Delta x \rightarrow 0$, we can use the Taylor's expansion series till 1st factor:









EVALUATION OF B_0



Noting that $B_1(0) = B_2(0) = B_0$. By replacing: $QB_{0} - K_{x}A\frac{dB_{1}}{dx}\Big|_{x=0} + Q_{w}B_{w} = QB_{0} + Q_{w}B_{0} - K_{x}A\frac{dB_{2}}{dx}\Big|_{x=0} = B_{0}\lambda_{1}e^{\lambda_{1}x}\Big|_{x=0} = B_{0}\lambda_{1}$ $\Rightarrow Q_{w}B_{w} - K_{x}AB_{0}\lambda_{1} = Q_{w}B_{0} - K_{x}AB_{0}\lambda_{2}$ $\frac{dB_{2}}{dx}\Big|_{x=0} = B_{0}\lambda_{2}e^{\lambda_{2}x}\Big|_{x=0} = B_{0}\lambda_{2}$ $Q_{w}B_{w} = Q_{w}B_{0} + K_{r}A(\lambda_{1} - \lambda_{2})B_{0} = B_{0}[Q_{w} + K_{r}A(\lambda_{1} - \lambda_{2})]$ $\lambda_1 - \lambda_2 = \frac{U}{2K_r} (1 + m_d - 1 + m_d) = \frac{U}{K_r} m_d$ $B_0 = B_w \frac{Q_w}{Q_w + K_x A \frac{U}{K_x} m_d}$ $B_0 = B_w \frac{Q_w}{Q_w + Q_w}$

It is worth noting that for pure advection $K_x = 0$, i.e. $m_d = 1$, the solution is the same determined in the standard S-P model:

$$B_0 = B_w \frac{Q_w}{Q_w + Q}$$









The eq. (1) in the most general form reads:

$$\frac{\mathrm{d}^2 \tilde{O}}{\mathrm{d}x^2} - \frac{U}{K_x} \frac{\mathrm{d}\tilde{O}}{\mathrm{d}x} - \frac{r_a}{K_x} \tilde{O} + \frac{r_d}{K_x} B = 0 \quad \longrightarrow \quad \frac{\mathrm{d}^2 y}{\mathrm{d}x^2} + a_1 \frac{\mathrm{d}y}{\mathrm{d}x} - a_0 y = c$$

The solution is the sum of the solution of the homogeneous eq. and a particular solution:

$$\frac{\mathrm{d}^2 y}{\mathrm{d}x^2} + a_1 \frac{\mathrm{d}y}{\mathrm{d}x} - a_0 y = 0 + \tilde{O}_p = k x^m e^{\lambda x}$$

The procedure to derive the solution of homegenous eq. is the same of the one developed for *BOD*:

$$\tilde{O} = c_1 e^{\lambda_1 x} + c_2 e^{\lambda_2 x} \qquad \text{with} \qquad \lambda_{1,2} = \frac{U}{2K_x} (1 \pm m_a)$$
$$m_a = \sqrt{1 + \frac{4r_a K_x}{U^2}} > 1$$

And the particular solution is:

$$\tilde{O}_p = \frac{r_d B_0}{r_a - r_d} e^{\frac{U}{2K_x}(1 \pm m_d)x}$$
 The particular integral changes when $x \ge 0$ or $x < 0$.









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Also in this case: $c_1 = 0$ for x > 0 $c_2 = 0$ for x < 0

The last Boundary Condition is $\tilde{O}(0) = \tilde{O}_0$, thus:

$$\begin{cases} \tilde{O}(0) = c_2 e^{\lambda_2 x} + \frac{r_d B_0}{r_a - r_d} e^{\frac{U}{2K_t} (1 - m_d) x} = \tilde{O}_0 \quad x \ge 0 \\ \tilde{O}(0) = c_1 e^{\lambda_1 x} + \frac{r_d B_0}{r_a - r_d} e^{\frac{U}{2K_t} (1 + m_d) x} = \tilde{O}_0 \quad x < 0 \end{cases}$$

$$\begin{cases} \tilde{O}_0 = c_2 + \frac{r_d B_0}{r_a - r_d} & x \ge 0 \\ \tilde{O}_0 = c_1 + \frac{r_d B_0}{r_a - r_d} & x < 0 \end{cases}$$

$$\longrightarrow c_1 = c_2 = \tilde{O}_0 - \tilde{O}_p = \tilde{O}_0 - \frac{r_d B_0}{r_a - r_d}$$







EVALUATION OF $\widetilde{m{O}}_{m{0}}$



As in the previous case \tilde{O}_0 is calculated by the following mass balance:

$$Q\tilde{O}_1(-\Delta x) - K_x A \frac{\mathrm{d}\tilde{O}_1}{\mathrm{d}x} \bigg|_{x=-\Delta x} = Q\tilde{O}_2(\Delta x) - K_x A \frac{\mathrm{d}\tilde{O}_2}{\mathrm{d}x} \bigg|_{x=\Delta x}$$

By expanding the Taylor's series around the inflow:

- $\tilde{O}_1(-\Delta x) \cong \tilde{O}_1(0)$
- $\tilde{O}_2(\Delta x) \cong \tilde{O}_2(0)$
- $\left. \frac{\mathrm{d}\tilde{O}_1}{\mathrm{d}x} \right|_{x=-\Delta x} \cong \left. \frac{\mathrm{d}\tilde{O}_1}{\mathrm{d}x} \right|_{x=0}$





By replacing and considering that $\tilde{O}_1(0) = \tilde{O}_2(0) = \tilde{O}_0$:







EVALUATION OF \widetilde{O}_0

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$$\left. \mathcal{Q}\tilde{\Theta}_{0} - K_{x}A\frac{\mathrm{d}\tilde{O}_{1}}{\mathrm{d}x} \right|_{x=0} = \mathcal{Q}\tilde{\Theta}_{0} + Q_{w}\tilde{O}_{0} - K_{x}A\frac{\mathrm{d}\tilde{O}_{2}}{\mathrm{d}x} \right|_{x=0}$$

Being generally $Q_{w}\tilde{O}_{0} \ll 1$: $\left. \frac{\mathrm{d}\tilde{O}_{1}}{\mathrm{d}x} \right|_{x=0} = \frac{\mathrm{d}\tilde{O}_{2}}{\mathrm{d}x} \right|_{x=0}$ The the

The two derivative are equal, thus there is no cusp.

To calculate the latter we derive the following expression:

$$\tilde{O}_{1,2} = c_{1,2}e^{\lambda_{1,2}x} + \frac{r_d B_0}{r_a - r_d}e^{\frac{U}{2K_x}(1 \pm m_d)x} \quad \text{with} \quad \lambda_{1,2} = \frac{U}{2K_x}(1 \pm m_a)$$

That is:

$$\frac{d\tilde{O}_1}{dx}\bigg|_{x=0} = c_1\lambda_1 + \frac{r_d B_0}{r_a - r_d} \frac{U}{2K_x} (1 + m_d) = c_2\lambda_2 + \frac{r_d B_0}{r_a - r_d} \frac{U}{2K_x} (1 - m_d) = \frac{d\tilde{O}_2}{dx}\bigg|_{x=0}$$

$$U_{x=0} = \frac{U_{x=0}}{V_{x=0}} \frac{V_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{V_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{V_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=0}} \frac{U_{x=0}}{V_{x=$$

$$c_1 \frac{0}{2K_x} (1+m_a) + \frac{r_d B_0}{r_a - r_d} \frac{0}{2K_x} (1+m_d) = c_2 \frac{0}{2K_x} (1-m_a) + \frac{r_d B_0}{r_a - r_d} \frac{0}{2K_x} (1-m_d)$$







EVALUATION OF \widetilde{O}_0



But
$$c_1 = c_2 = c$$
:
 $\mathcal{Z}m_a$
 $c(1 + m_a - 1 + m_a) = \frac{r_d B_0}{r_a - r_d} (1 - m_d - 1 - m_d) \longrightarrow c_{1,2} = -\frac{r_d B_0}{r_a - r_d} \frac{m_d}{m_a}$

By replacing the coefficients into the expression of \tilde{O}_0 , finally we find:





EXERCISE: \widetilde{O} IN ESTUARY



Let's study the *BOD* and *DO* distribution of the following scenario.

<u>Data:</u>

- $Q = 30 \text{ m}^3/\text{s}$
- $Q_w = 5 \text{ m}^3/\text{s}$
- $B_w = 50 \text{ mg/l}$
- $U = 0.15 \text{ m/s} \longrightarrow U = 12.96 \text{ km/d}$
- $K_{\chi} = 380 \text{ m}^2/\text{s} \longrightarrow K_{\chi} = 32.83 \text{ km}^2/\text{d}$
- $r_d = 0.5 \text{ d}^{-1}$
- $r_a = 0.9 \text{ d}^{-1}$
- $O_s = 7.5 \text{ mg/l}$



The regime of *O* is determined by the values of the following parameter:

$$\frac{r_d K_x}{U^2} = \frac{0.5 \cdot 32.83}{12.96^2} = 0.098 \qquad \qquad \frac{r_a K_x}{U^2} = \frac{0.9 \cdot 32.83}{12.96^2} = 0.176$$

The two parameters are within the range [0.05, 20], thus we have to take into account both <u>advection</u> and <u>dispersion</u>.







EXERCISE: *O* IN ESTUARY



The equations describing the dynamics of B and \tilde{O} depend on the parameter m_d and m_a :

$$m_d = \sqrt{1 + \frac{4r_d K_x}{U^2}} = \sqrt{1 + 4 \cdot 0.098} = 1.18$$

$$m_a = \sqrt{1 + \frac{4r_a K_x}{U^2}} = \sqrt{1 + 4 \cdot 0.176} = 1.31$$

The *BOD* distribution is given by:

$$B = B_0 e^{\frac{U}{2K_x}(1-m_d)x} \quad x \ge 0$$
$$B = B_0 e^{\frac{U}{2K_x}(1+m_d)x} \quad x < 0$$

where the value of B_0 is equal to:

$$B_0 = B_w \frac{Q_w}{Q_w + Qm_d} = 80 \frac{5}{5 + 30 \cdot 1.18} = 9.9 \text{ mg/l}$$

By replacing all the parameters into the system, we find:

$$\begin{cases} B^+ = 9.9 \cdot e^{-0.035 \cdot x} \\ B^- = 9.9 \cdot e^{+0.430 \cdot x} \end{cases}$$









Similarly the \tilde{O} distribution is described by the following system:

$$\begin{cases} \tilde{O}^{+} = \frac{r_d B_0}{r_a - r_d} \left[e^{\frac{U}{2K_x} (1 - m_d)x} - \frac{m_d}{m_a} e^{\frac{U}{2K_x} (1 - m_a)x} \right] = 12.38 [e^{-0.035 \cdot x} - 0.903 e^{-0.060 \cdot x}] \\ \tilde{O}^{-} = \frac{r_d B_0}{r_a - r_d} \left[e^{\frac{U}{2K_x} (1 + m_d)x} - \frac{m_d}{m_a} e^{\frac{U}{2K_x} (1 + m_a)x} \right] = 12.38 [e^{+0.430 \cdot x} - 0.903 e^{+0.455 \cdot x}] \end{cases}$$

Finally, the *DO* distribution can be trivially determined by: $O = O_s - \tilde{O}$

The numerical analysis of \tilde{O} shows that the critic condition is:

$$x_c = 17.3 \text{ km}$$

 $\tilde{O}_c = 2.8 \text{ mg/l} \longrightarrow O_c = 7.5 - 2.8 = 4.7 > 4.0 \text{ mg/l}$

The critic condition shows accettable concentration of *DO*.







EXERCISE: *O* IN ESTUARY









