



Università degli Studi di Padova

LESSON 2: MOLECULAR DIFFUSION







<u>Hp</u>: - Fluid at rest: u(x, t) = 0

- Solute mass conservation:
$$\frac{\mathrm{d}M}{\mathrm{d}t} = \mathrm{const}$$

Fick stated:

• The mass flux of solute depends on its concentration gradient.

 $\boldsymbol{q}=f(\boldsymbol{\nabla} c)$

- The process is <u>isotropic</u> (equal characteristics along all directions).
- The process is <u>homogeneous</u> (it does not depend on space and position).
- The process is linear.

$$q = -D \nabla c$$

Definition of Gradient ∇ Scalar: $\nabla c = \begin{bmatrix} \frac{\partial c}{\partial x} & \frac{\partial c}{\partial y} & \frac{\partial c}{\partial z} \end{bmatrix}$ Vector: $\nabla a = \begin{bmatrix} \frac{\partial a_1}{\partial x_1} & \frac{\partial a_1}{\partial x_2} & \frac{\partial a_1}{\partial x_3} \\ \frac{\partial a_2}{\partial x_1} & \frac{\partial a_2}{\partial x_2} & \frac{\partial a_2}{\partial x_3} \\ \frac{\partial a_3}{\partial x_1} & \frac{\partial a_3}{\partial x_2} & \frac{\partial a_3}{\partial x_3} \end{bmatrix}$







FICK LAW



D is the molecular diffusion coefficient or diffusivity.

$$D = D(T, p)$$
 or $D = D(T, p, c_s)$ $[m^2/s]$
 T is the fluid Temperature \longrightarrow $D \propto T$
 p is the fluid pressure \longrightarrow $D \propto 1/p$

Values of diffusion coefficients (liquid)		
Species pair (solute – solvent)	Temperature (°C)	D (cm²/s)
Acetone (dis) – water (I)	25	1.16×10⁵
Air (dis) – water (I)	25	2.00×10-5
Ammonia (dis) – water (l)	25	1.64×10⁵
Argon (dis) – water (I)	25	2.00×10-5
Benzene (dis) – water (l)	25	1.02×10⁻⁵
Bromine (dis) – water (I)	25	1.18×10⁻⁵
Carbon monoxide (dis) – water (l)	25	2.03×10-⁵
Carbon dioxide (dis) – water (I)	25	1.92×10⁻⁵
Chlorine (dis) – water (l)	25	1.25×10⁵
Ethane (dis) – water (l)	25	1.20×10⁻⁵
Ethanol (dis) – water (I)	25	0.84×10⁻⁵
Ethylene (dis) – water (I)	25	1.87×10⁻⁵
Helium (dis) – water (I)	25	6.28×10⁻⁵
Hydrogen (dis) – water (l)	25	4.50×10⁻⁵









A-D equation is based on the mass conservation

 $\begin{cases} M = \int_{\forall} c_s \, \mathrm{d} \forall \\ \frac{\mathrm{d}M}{\mathrm{d}t} = 0 \end{cases}$

∀ is an arbitrary Material Volume

By replacing the first equation into the second one:

$$\frac{\mathrm{d}M}{\mathrm{d}t} = \frac{\mathrm{d}}{\mathrm{d}t} \int_{\forall} c_s \, \mathrm{d}\forall = 0$$

$$Transport th:: \quad \frac{\mathrm{d}\Psi}{\mathrm{d}t} = \int_{\forall} \left(\frac{\mathrm{d}\psi}{\mathrm{d}t} + \psi \,\nabla \cdot \boldsymbol{u}\right) \, \mathrm{d}\forall$$

$$\frac{\mathrm{d}M}{\mathrm{d}t} = \int_{\forall} \left(\frac{\mathrm{d}c_s}{\mathrm{d}t} + c_s \nabla \cdot \boldsymbol{u}_s\right) \, \mathrm{d}\forall = 0$$

Definition of velocity divergence ∇ \cdot :

$$\nabla \cdot \boldsymbol{u} = \frac{\partial u_x}{\partial x} + \frac{\partial u_y}{\partial y} + \frac{\partial u_z}{\partial z}$$









The integral is zero for any \forall , that is:

$$\frac{\mathrm{d}c_s}{\mathrm{d}t} + c_s \nabla \cdot \boldsymbol{u}_s = 0$$
$$\frac{\partial c_s}{\partial t} + \boldsymbol{u}_s \cdot \nabla c_s + c_s \nabla \cdot \boldsymbol{u}_s = 0$$
$$\frac{\partial c_s}{\partial t} + \nabla \cdot (c_s \, \boldsymbol{u}_s) = 0$$

If we rearreange the expression by decomposing the absolute velocity of solute, we obtain:

$$\frac{\partial c_s}{\partial t} + \nabla \cdot [c_s(\mathbf{U} + \mathbf{v}_s)] = 0$$

$$\frac{\partial c_s}{\partial t} + \nabla \cdot (c_s \mathbf{U}) = -\nabla \cdot (c_s \mathbf{v}_s)$$
Differential equation of the solute mass balance

7 unknown parameters VS 5 equations (1 solute mass balance, 1 continuity eq., 3 momentum eqs.)









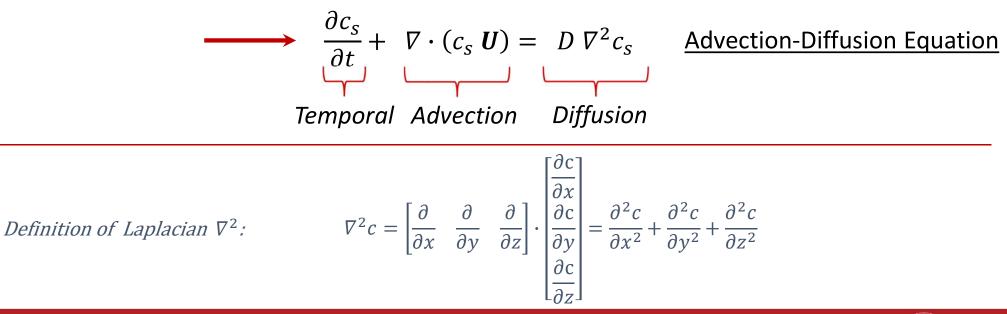
The parameters can be reduced through the Fick law. Indeed:

$$\boldsymbol{q_s^r} = c_s \, \boldsymbol{\nu_s} = -D \, \nabla c_s$$

And then:

 $-\nabla \cdot (c_s \, \boldsymbol{v}_s) = \nabla \cdot (D \, \nabla c_s)$ $-\nabla \cdot (c_s \, \boldsymbol{v}_s) = D \, \nabla \cdot \nabla c_s$ $-\nabla \cdot (c_s \, \boldsymbol{v}_s) = D \, \nabla^2 c_s$

By replacing the latter into solute mass balance:











Let's consider the mass balance equation $c_s = \rho$:

$$\frac{\mathrm{d}M}{\mathrm{d}t} = \int\limits_{\forall} \left(\frac{\mathrm{d}\rho}{\mathrm{d}t} + \rho \nabla \cdot \boldsymbol{u}_{\boldsymbol{s}} \right) \, \mathrm{d} \forall = 0$$

Being the Control Volume ∀ arbitrary, the integrating function must be:

$$\frac{\partial \rho}{\partial t} + \boldsymbol{u}_{\boldsymbol{s}} \cdot \nabla \rho + \rho \nabla \cdot \boldsymbol{u}_{\boldsymbol{s}} = 0$$

Usually water is assumed to be <u>uncompressible</u> and <u>homogeneous</u>. It means: $\rho = \rho(\mathbf{x}, t) = const$ Then:

$$\frac{\partial \rho}{\partial t} + \boldsymbol{u} \cdot \nabla \rho + \rho \nabla \cdot \boldsymbol{u} = 0$$

$$\longrightarrow \quad \nabla \cdot \boldsymbol{u} = 0 \qquad \text{3D C}$$

<u>3D Continuity Equation</u>







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By replacing continuity equation into advection-diffusion equation, we have:

$$\frac{\partial c_s}{\partial t} + \boldsymbol{U} \cdot \nabla c_s + c_s \nabla \cdot \boldsymbol{U} = D \nabla^2 c_s$$

And finally:

$$\frac{\partial c_s}{\partial t} + \boldsymbol{U} \cdot \nabla c_s = D \nabla^2 c_s$$

$$\frac{\text{Advection-Diffusion Equation}}{\text{for incompressible fluid}}$$

If we consider fluid at rest, i.e. U = 0, the latter equation is simplified in the following:

$$\frac{\partial c_s}{\partial t} = D \nabla^2 c_s$$
 Pure Diffusion Equation









Let's start from the A-D equation:

$$\frac{\partial c}{\partial t} + \boldsymbol{U} \cdot \boldsymbol{\nabla} C = D \, \boldsymbol{\nabla}^2 C$$

<u>Hp</u>: - non-reactive solute: dM/dt = 0- homogeneous and uncompressible fluid: $\rho = 0$

When the fluid is at rest, i.e. *U*=0, then:

$$\frac{\partial c_s}{\partial t} = D \nabla^2 c_s$$
 Diffusion Equation

This equation is <u>linear</u> and it can be solved analitycally!

Why do we solve molecular diffusion when this phenomenon is negligible at large scale?

Analitycal solution of molecular diffusivity and turbulent diffusion is the same, hence this solution can be extended to the turbulent case (under certain conditions).









To solve the Diffusion equation the <u>initial</u> and the <u>boundary conditions</u> have to be correctly set.

Initial condition: $c(\mathbf{x},t) \rightarrow c(\mathbf{x},0) = c_0(\mathbf{x})$ Initial distribution of solute concentration

Boundary condition: $c(\mathbf{x}, t) \rightarrow c(\mathbf{x}_0, t)$

Evolution with t of c in some sections of the domain

Three kinds of BCs exist:

- Dirichlet condition: c = c(t) in x_0
- Neumann condition: ∇c in $x_0 \rightarrow q^r = c v_s$ in x_0
- Cauchy condition: Dirichlet condition + Neumann Condition

N.B. In our applications also BCs are linear!









Why is it fundamental?

Since the problem is linear, one can use the superimposition of the effects to solve complex problem.

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complex problem = \sum elementary problems
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→ We will know their solution

Let's consider the following **1-D Diffusion problem**:

$$\begin{cases} \frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2} \\ c(x,0) = M\delta(x) & \text{usually for } x = 0 \\ c(x,t) = 0 & \text{for } x \to \pm \infty \end{cases} \xrightarrow{\text{Initial condition}} \\ \frac{\text{Boundary conditions}}{\text{Boundary conditions}} \end{cases}$$

In this case the concentration *c* is defined as:

$$c = \lim_{\Delta x \to 0} \frac{\Delta M}{\Delta x} \qquad [kg/m]$$









Initial condition $c(x, 0) = M \delta(x)$ is the <u>injection point</u>.

Boundary conditions c(x, t) = 0 for $x \to \pm \infty$ means domain <u>undefinetely extended</u>. Noting that:

- i) Actually, c = c(x, t, M, D)ii) Due to linearity, $c \propto M$ c = M f(x, t, D)
- iii) For the mass conservation, $M = \int_{-\infty}^{+\infty} c(x, t) dx$
- iv) For isotropy the concentration distribution is symmetrical

 $\delta(x)$ is Dirac Delta Function, that is defined as following:

Dirac Delta Function has the property: f(

s following:
$$\int_{-\infty} \delta(x) \, dx = 1$$
$$\int_{-\infty}^{\infty} f(x)\delta(x - x_0) \, dx = \begin{cases} 0 & x \neq x_0 \\ f(x_0) & x = x_0 \end{cases}$$

 $c + \infty$

ICEA/

Environmental Fluid Mechanics – Lesson 2: Molecular Diffusion



 $\delta(x)$

 \rightarrow max