

# MgB<sub>2</sub>: a two-gap superconductor for practical application

*M. Tropeano*



*In collaboration with*



*CNR-SPIN, Genoa, Italy*



*Dipartimento di Fisica, University of Genoa, Italy*

## Introduction:

- Chronology of  $T_c$  until 2001
- Discovery of SC at 39K in  $MgB_2$
- Relevant results in the firsts 5 months after the discovery

## Properties of $MgB_2$ :

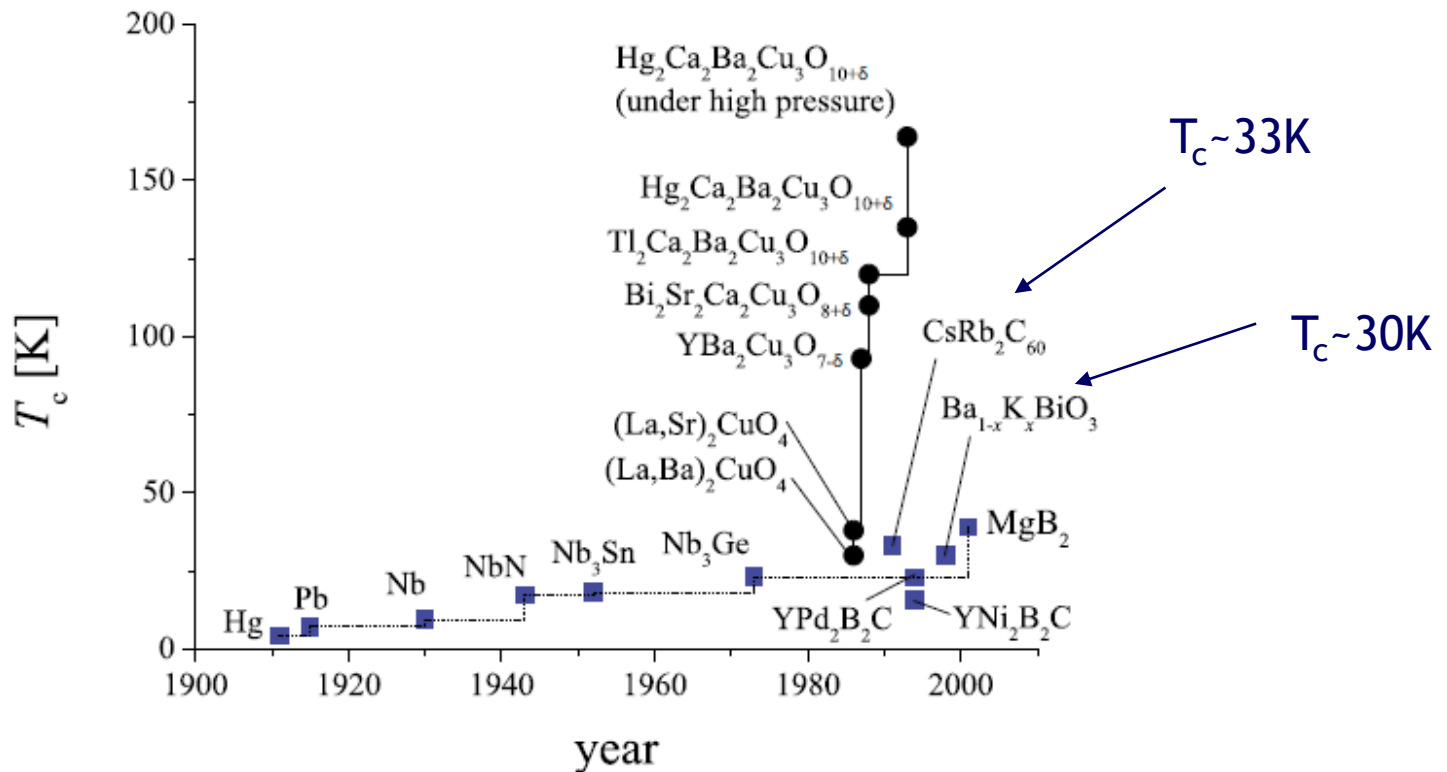
- Crystal and electronic structure
- Phonon modes, e-ph coupling, isotope effect
- Exp evidence of two gaps
- Impurities and irradiation effect
- Resistivity of  $MgB_2$  and its correlation with  $T_c$

In the first part the discovery of  $MgB_2$  and the relevant results in the firsts months will be briefly discussed to show how  $MgB_2$  since the beginning has been considered a material suitable for applications.

The principal aspects of the physics of  $MgB_2$  will be also described: discussion could not be exhaustive, the aim is to introduce the fundamental properties showing the peculiar characteristics hoping to fill you with curiosity about this compound

## Applications of MgB

- General properties of SC cables available on the market
- PIT fabrication options: in-situ and ex-situ
- Effect of deformation and thermal treatment on
- Powder and conductor optimization
- Mechanical properties
- Requirements for applications
- Demonstrators of MgB<sub>2</sub> technology
- Market potential



1911: zero resistivity at about 4K in Hg followed SC in other pure element

1940-1950: A15 compounds with maximum T<sub>c</sub> at about 23K in Nb<sub>3</sub>Sn

BCS theory  
framework

1986: Cu oxide SC with a maximum T<sub>c</sub> at 138K at ambient pressure

→ New theoretical  
interpretation

How much will T<sub>c</sub> increase in non Cu-oxyde superconductors?

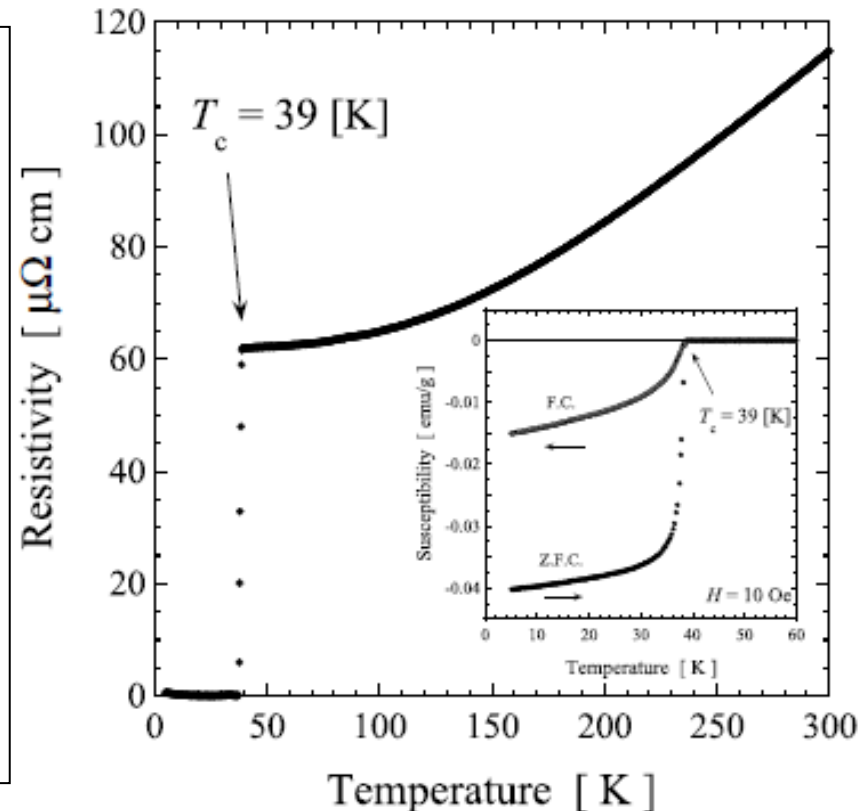
NATURE | VOL 410 | 1 MARCH 2001 | www.nature.com

# Superconductivity at 39 K in magnesium diboride

Jun Nagamatsu\*, Norimasa Nakagawa\*, Takahiro Muranaka\*,  
Yuji Zenitani\* & Jun Akimitsu\*†

\* Department of Physics, Aoyama-Gakuin University, Chitosedai, Setagaya-ku,  
Tokyo 157-8572, Japan

† CREST, Japan Science and Technology Corporation, Kawaguchi, Saitama 332-  
0012, Japan



MgB<sub>2</sub> has opened a new frontier into the physical properties and application of SC.

The limit of  $T_c$  in metallic superconductors had been believed to be about 30 K in the framework of the BCS theory and this discovery of unexpected high- $T_c$  in this simple binary intermetallic compound has triggered enormous interests in the world.

If you run across a new metal, or an old one, **cool it down!**  
You might get a pleasant surprise (Paul M. Grant).

NATURE | VOL 410 | 8 MARCH 2001 | www.nature.com

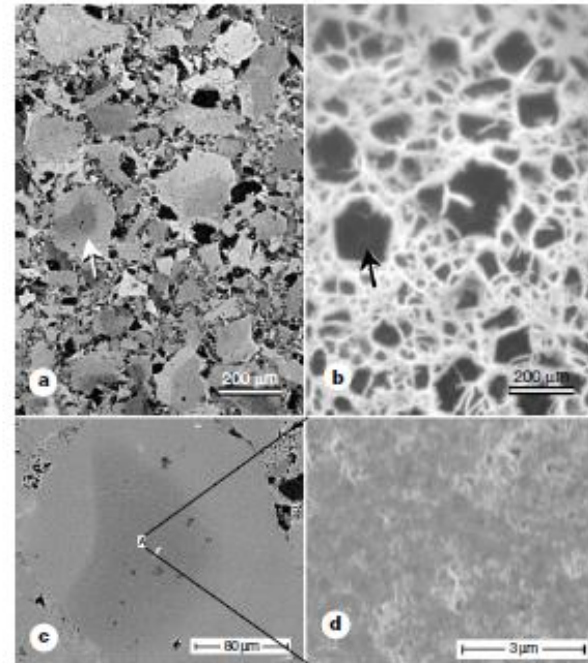
## Strongly linked current flow in polycrystalline forms of the superconductor $MgB_2$

D. C. Larbalestier<sup>\*†</sup>, L. D. Cooley<sup>\*</sup>, M. O. Rikel<sup>\*</sup>, A. A. Polyanskii<sup>\*</sup>, J. Jiang<sup>\*</sup>, S. Patnaik<sup>\*</sup>, X. Y. Cai<sup>\*</sup>, D. M. Feldmann<sup>\*</sup>, A. Gurevich<sup>\*</sup>, A. A. Squitieri<sup>\*</sup>, M. T. Naus<sup>\*</sup>, C. B. Eom<sup>\*†</sup>, E. E. Hellstrom<sup>\*†</sup>, R. J. Cava<sup>‡</sup>, K. A. Regan<sup>‡</sup>, N. Rogado<sup>‡</sup>, M. A. Hayward<sup>‡</sup>, T. He<sup>‡</sup>, J. S. Slusky<sup>‡</sup>, P. Khalifah<sup>‡</sup>, K. Inumaru<sup>‡</sup> & M. Haas<sup>‡</sup>

<sup>\*</sup> Applied Superconductivity Center, University of Wisconsin–Madison, 1500 Engineering Drive, Madison, Wisconsin 53706, USA

<sup>†</sup> Department of Materials Science and Engineering, University of Wisconsin–Madison, 1509 University Avenue, Madison, Wisconsin 53706, USA

<sup>‡</sup> Department of Chemistry and Princeton Materials Institute, Princeton University, Princeton, New Jersey 08544, USA



**Cu-oxide SC:** high  $T_c$  but large anisotropy and weak coupling across grain boundaries

→ Transport  $J_c$  in untextured samples low and sensitive to magnetic field

**$MgB_2$ :** despite the multiphase, untextured, microscale subdivided nature of the first samples supercurrents flow in the material without exhibiting strong sensitivity to weak magnetic field

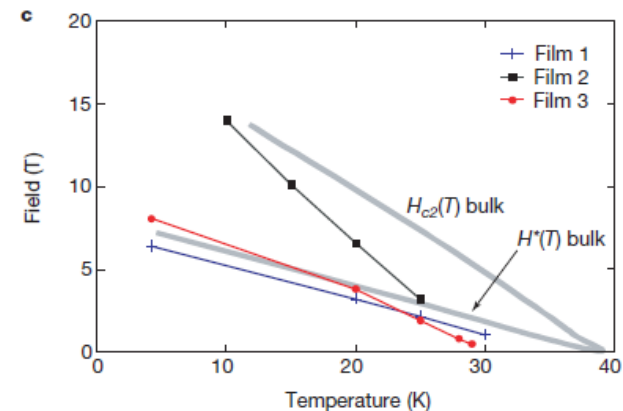
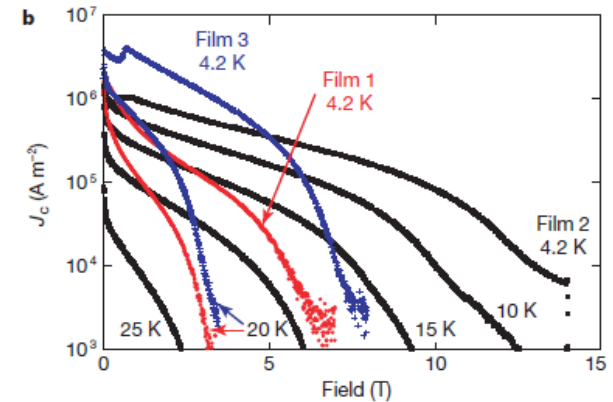
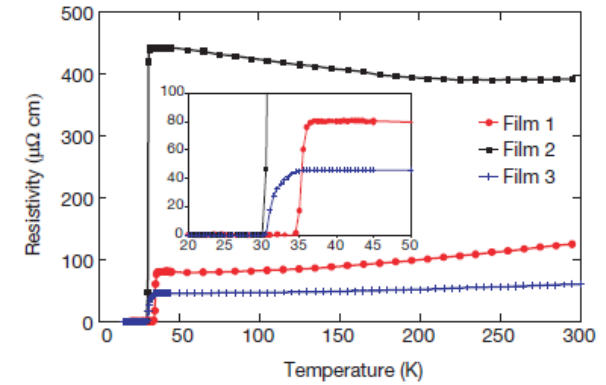
Remarkable absence of “weak link”, no need of high degree of texturing

NATURE | VOL 411 | 31 MAY 2001 | www.nature.com

**High critical current density and enhanced irreversibility field in superconducting MgB<sub>2</sub> thin films**

C. B. Eom<sup>\*†</sup>, M. K. Lee<sup>\*</sup>, J. H. Choi<sup>\*†</sup>, L. J. Belenky<sup>\*</sup>, X. Song<sup>†</sup>, L. D. Cooley<sup>†</sup>, M. T. Naus<sup>†</sup>, S. Patnaik<sup>†</sup>, J. Jiang<sup>†</sup>, M. Rikel<sup>†</sup>, A. Polyanskii<sup>†</sup>, A. Gurevich<sup>†</sup>, X. Y. Cai<sup>†</sup>, S. D. Bu<sup>\*</sup>, S. E. Babcock<sup>\*†</sup>, E. E. Hellstrom<sup>\*†</sup>, D. C. Larbalestier<sup>\*†</sup>, N. Rogado<sup>‡</sup>, K. A. Regan<sup>‡</sup>, M. A. Hayward<sup>‡</sup>, T. He<sup>‡</sup>, J. S. Slusky<sup>‡</sup>, K. Inumaru<sup>‡</sup>, M. K. Haas<sup>‡</sup> & R. J. Cava<sup>‡</sup>

<sup>\*</sup> Department of Materials Science and Engineering, 1509 University Avenue; and  
<sup>†</sup> Applied Superconductivity Center, 1500 Engineering Drive, University of Wisconsin, Madison, Wisconsin 53706, USA  
<sup>‡</sup> Department of Chemistry and Princeton Materials Institute, Princeton University, Princeton, New Jersey 08544, USA



3 thin films:

PLD and different annealing with unintentional substitution of O on B

**Enhanced J<sub>c</sub> behaviour** in magnetic field (10<sup>6</sup> A/cm<sup>2</sup> at 1T and 10<sup>5</sup> at 10T) and **enhanced H<sub>irr</sub>** (greater than 14T at 4.2K) respect to bulk

**MgB<sub>2</sub> has potential for in field application**

Performance of MgB<sub>2</sub> can rival and eventually surpass that of existing SC wires

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## Enhancement of the high-magnetic-field critical current density of superconducting MgB<sub>2</sub> by proton irradiation

Y. Bugoslavsky<sup>\*†</sup>, L. F. Cohen<sup>\*</sup>, G. K. Perkins<sup>\*</sup>, M. Polichetti<sup>‡</sup>,  
T. J. Tate<sup>\*</sup>, R. Gwilliam<sup>§</sup> & A. D. Caplin<sup>\*</sup>

<sup>\*</sup> Centre for High Temperature Superconductivity, Blackett Laboratory, Imperial College, London SW7 2BZ, UK

<sup>†</sup> General Physics Institute, 117942, Vavilov st38, Moscow, Russia

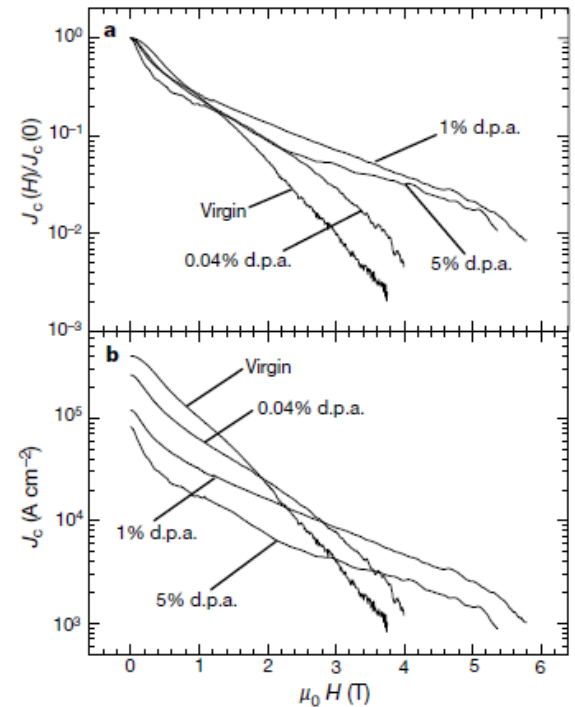
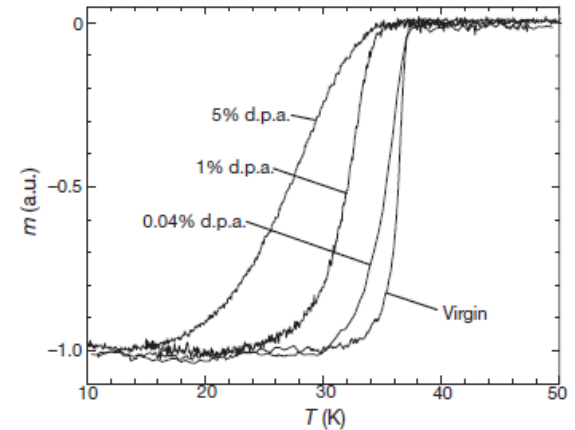
<sup>‡</sup> INFN–Dipartimento di Fisica, Università di Salerno, via S. Allende, Baronissi, I-84081 (Salerno), Italy

<sup>§</sup> EPSRC Ion Beam Centre, University of Surrey, Guildford, Surrey GU2 7RX, UK

Magnitude and field dependence of  $J_c$  are related to defect that can “pin” magnetic vortices.

A modest level of atomic disorder induced by proton irradiation enhance the pinning of vortices with a significant oncrease of  $J_c$  at high field

Chemical doping or mechanical processing can generate similar levels of disorder improving the performance



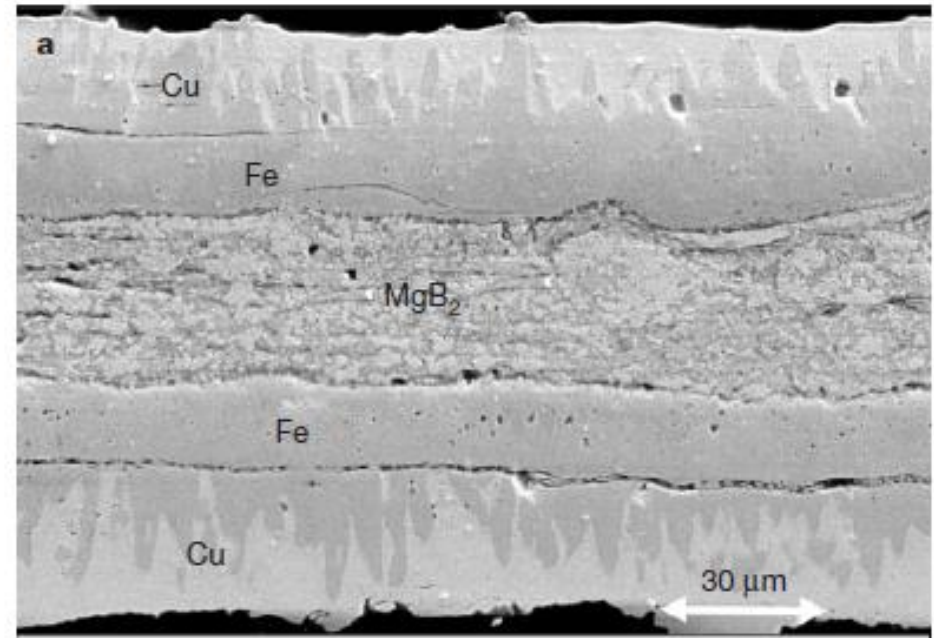


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## High critical currents in iron-clad superconducting MgB<sub>2</sub> wires

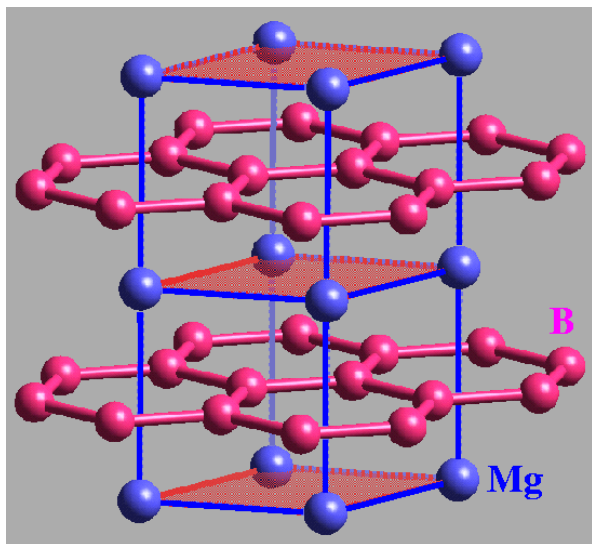
S. Jin, H. Mavoori, C. Bower & R. B. van Dover

Agere Systems/Lucent Technologies, Murray Hill, New Jersey 07974, USA



Technically useful bulk superconductors must have high  $J_c$ , at operating temperatures. They also require a normal metal cladding to provide parallel electrical conduction, thermal stabilization, and mechanical protection of the generally brittle superconductor cores.

The paper demonstrated a transport  $J_c$  in excess of 85,000 A/cm<sup>2</sup> at 4.2 K (a value high enough for power transmission cable)



MgB<sub>2</sub> has the AlB<sub>2</sub> structure with  $a=3.08\text{\AA}$  and  $c=3.51\text{\AA}$  (space group P6/mmm)

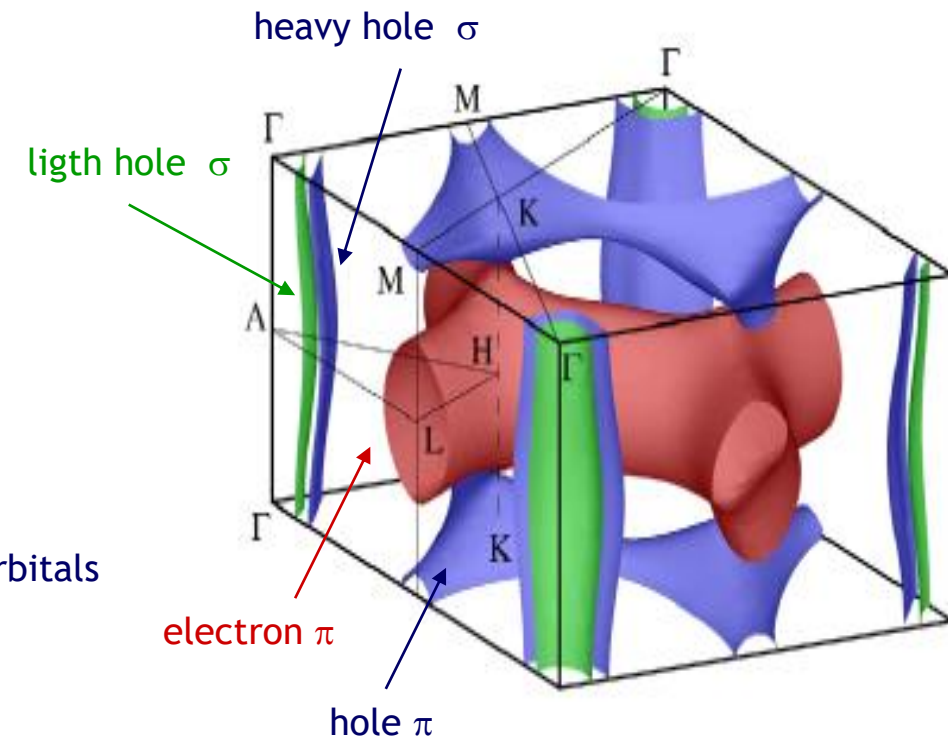
- stacks of 2D honeycomb B layer
- Mg atoms between the B layer at the centre of the hexagons

Band structure is derived by B electrons.

There are 4 conduction bands formed by

- two  $\sigma$  bands derived from the  $\sigma$ -bonding  $p_{x,y}$  orbitals
- two  $\pi$  bands derived by  $\pi$ -bonding (hole like) and anti-bonding (electron like)  $p_z$  orbitals

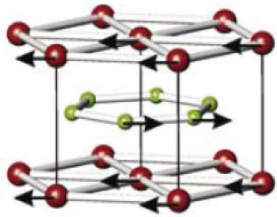
$\sigma$  states are 2D, confined in the B planes  
 $\pi$  states are 3D extending in all directions



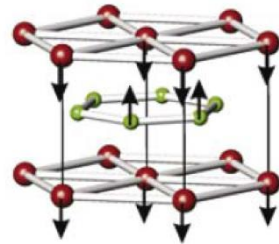
**Small Impurity interband scattering**

In plane and out of plane orbitals, different parity

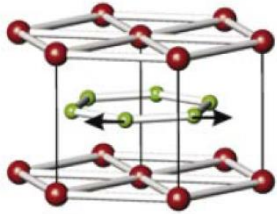
$E_{1u}$



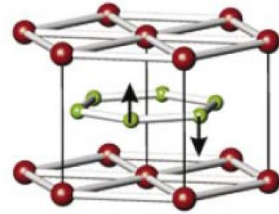
$A_{2u}$



$E_{2g}$



$B_{1g}$



The investigation of phonons is important because they play a crucial role in superconductivity

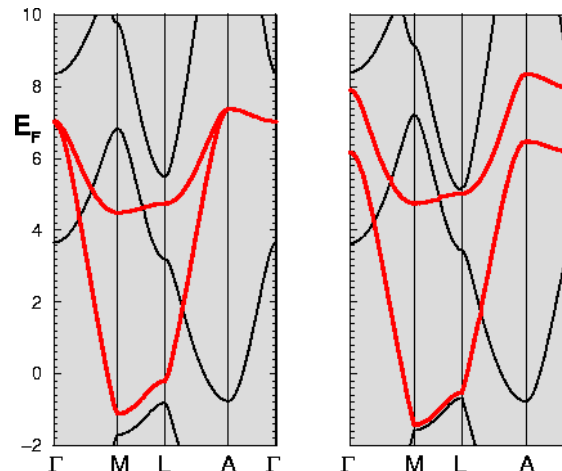
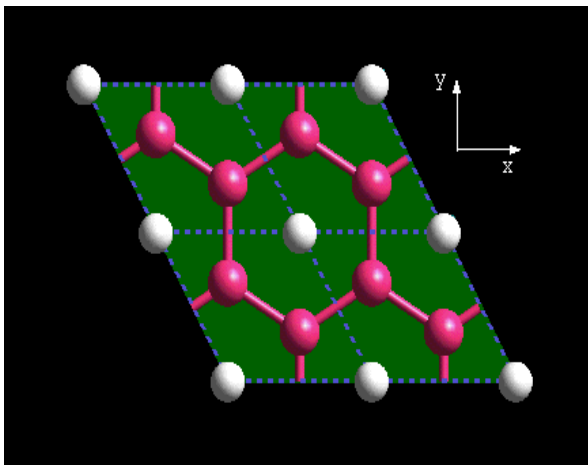
$E_{1u}$ : in plane B and Mg displacements

$A_{2u}$ : out of plane B and Mg displacements

$E_{2g}$ : B in plane bond stretching

$B_{1g}$ : B out of plane bond bending

Theoretical investigation show planar boron  $\sigma$  bands coupled with  $E_{2g}$  phonon modes



B motions change the boron orbital overlap with shift and splitting near the Fermi level



**Strong EP coupling**

In order to understand the role of phonons in SC the isotope effect is a crucial information

BCS isotope coefficient

$$\alpha = - \frac{\Delta \ln(T_c)}{\Delta \ln M}$$

Critical temperature

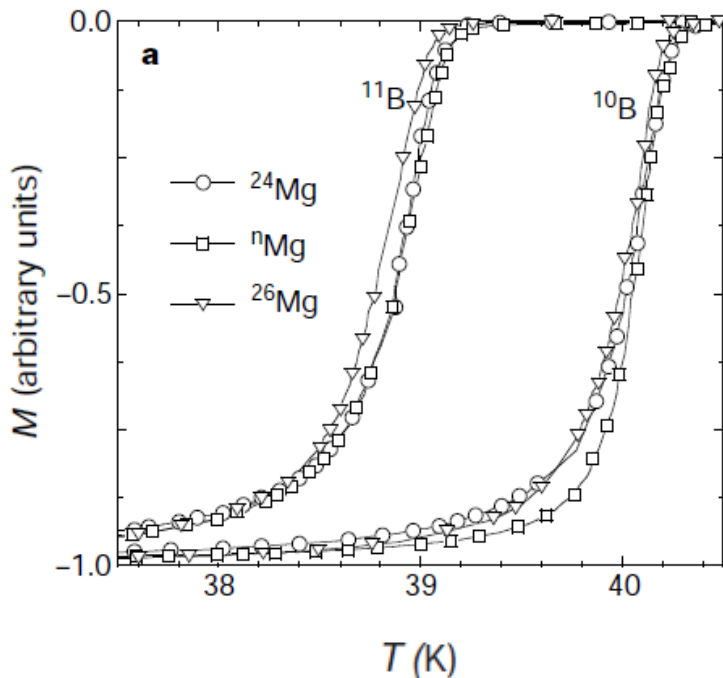
Isotopic mass of the element

Multi-element materials

$$\alpha_{tot} = \sum_i \alpha_i$$

MgB<sub>2</sub>

$$\alpha_{MgB_2} = \alpha_B + \alpha_{Mg}$$



The measurement of the isotope effect as the first evidence of the phonon mediated character of the coupling mechanism

$$\triangleright T_c(^{24}\text{MgB}_2) - T_c(^{26}\text{MgB}_2) \approx 0.07 \text{ K} \longrightarrow a_{\text{Mg}} \approx 0.02$$

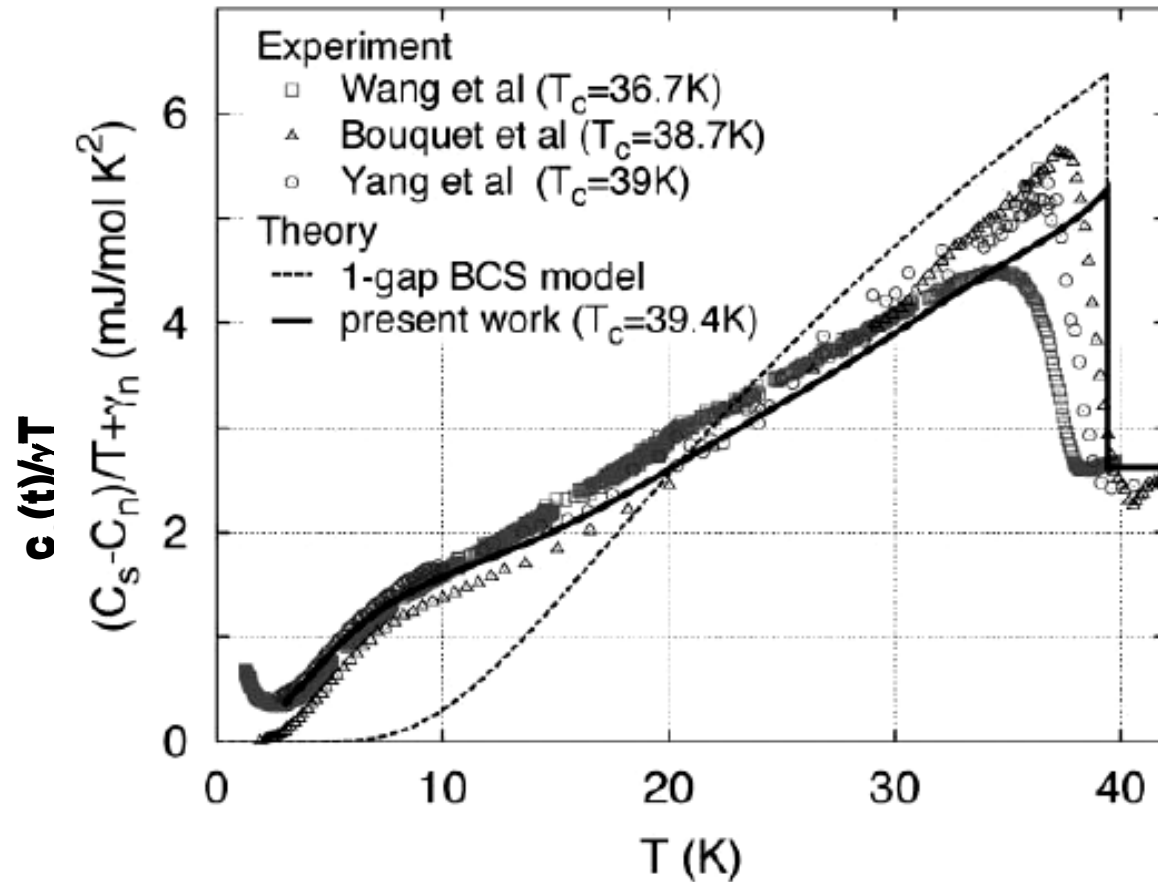
$$\triangleright T_c(\text{Mg}^{10}\text{B}_2) - T_c(\text{Mg}^{11}\text{B}_2) \approx 1 \text{ K} \longrightarrow a_B \approx 0.3$$

$$\alpha_{tot} = \alpha_B + \alpha_{Mg} \approx 0.32$$

$$\alpha_{\text{BCS}} = 0.5$$

$$\alpha(\text{MgB}_2) < a_{\text{BCS}} \quad ?$$

The anomalous value of the isotope coefficient is still an open question



$$T_c = 38.5 \text{ K}$$

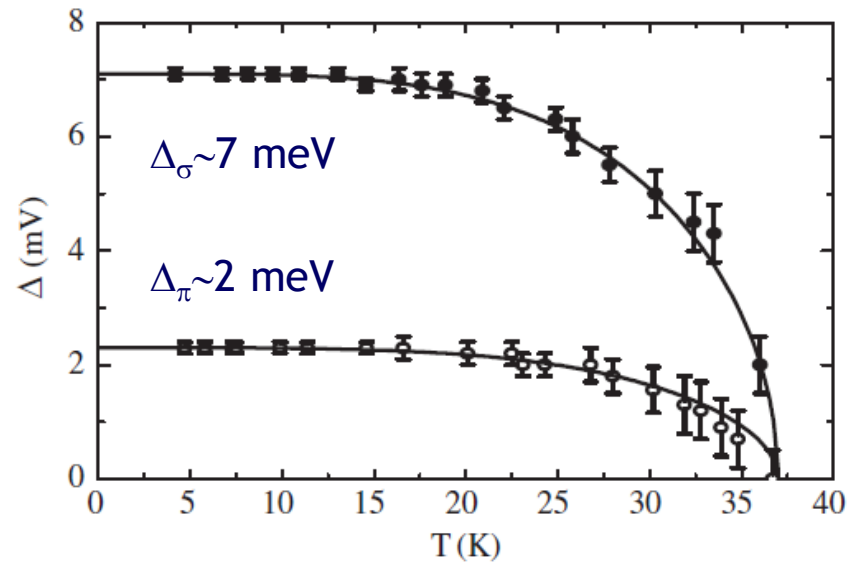
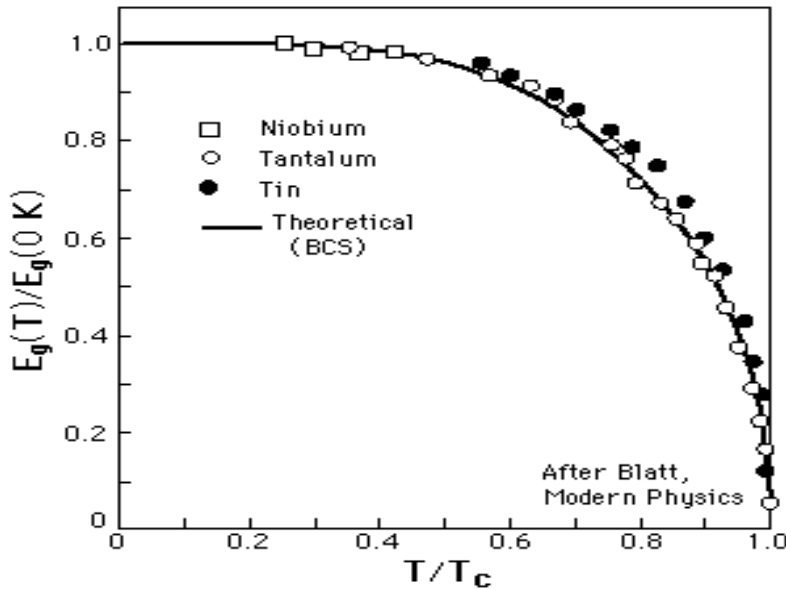
$$\Delta_\sigma(0) = 6.3 \text{ meV}$$

$$\Delta_\pi(0) = 1.8 \text{ meV}$$

Electronic specific heat data cannot be explained by the one gap BCS model with  $T_c$  of 39.4K. At lower T data values are higher than BCS showing a decrease to zero suggestive of another transition, a strong indication of a second, smaller gap.

The two gaps model is in excellent agreement with the experiments,

MgB<sub>2</sub> is a conventional BCS superconductor: Cooper pairs form via e-ph interaction.  
But MgB<sub>2</sub> is not an ordinary superconductor with a one temperature dependent energy gap  $\Delta(T)$



Two gaps exist in MgB<sub>2</sub> (data in figure are from Scanning tunneling spectroscopy) and disappear at  $T_c$ .

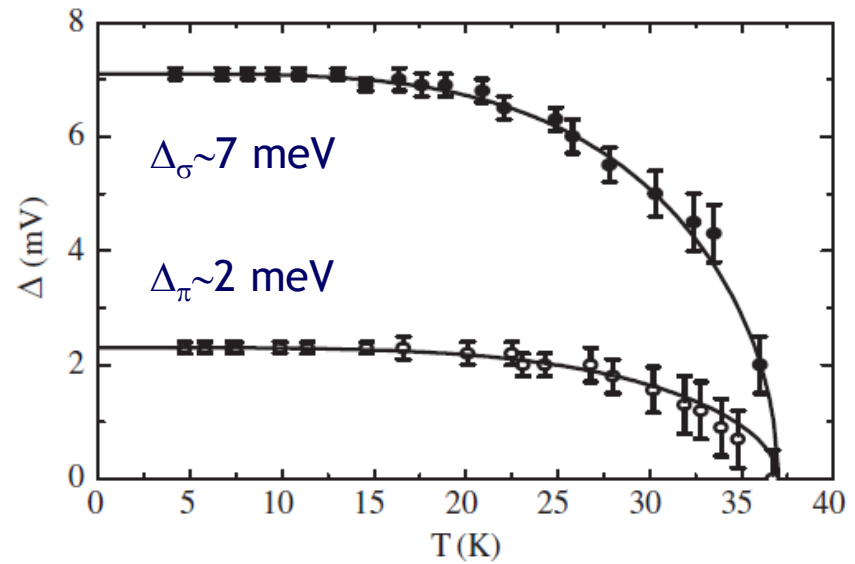
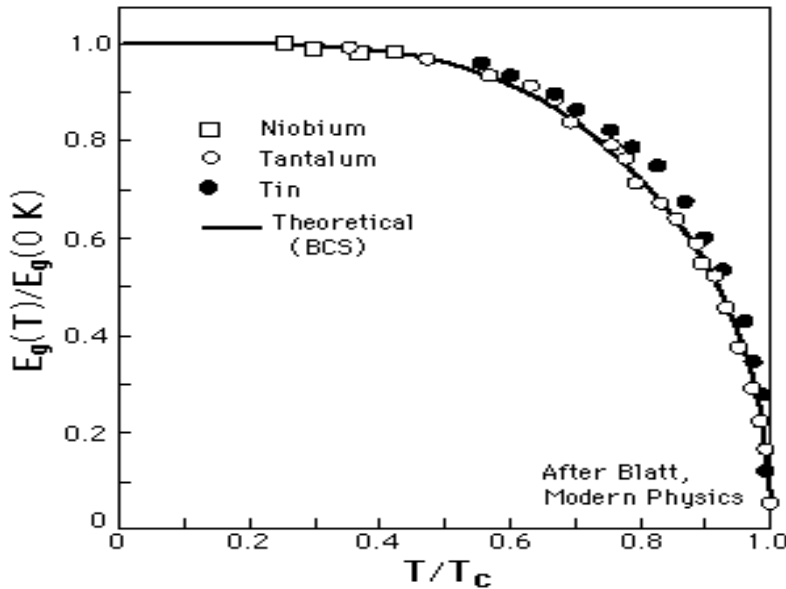
Two gaps arise from the existence of the two bands, but the small interband scattering distinguishes MgB<sub>2</sub> from other SC's with multiple bands at FS where the interband scattering smears out the distinct characteristic of the bands

Small Impurity interband scattering



Cooper pairs are not mixed maintaining their characteristics with two perfectly separated gaps

MgB<sub>2</sub> is a conventional BCS superconductor: Cooper pairs form via e-ph interaction.  
But MgB<sub>2</sub> is not an ordinary superconductor with a one temperature dependent energy gap  $\Delta(T)$



BCS equation for  $T_c$

$$T_c = \frac{\omega_D}{1.2} \exp\left(-\frac{1}{\lambda_{eff}}\right)$$

$\lambda_{\sigma\sigma}=0.96$	$\lambda_{\sigma\pi}=0.17$
$\lambda_{\pi\sigma}=0.23$	$\lambda_{\pi\pi}=0.29$

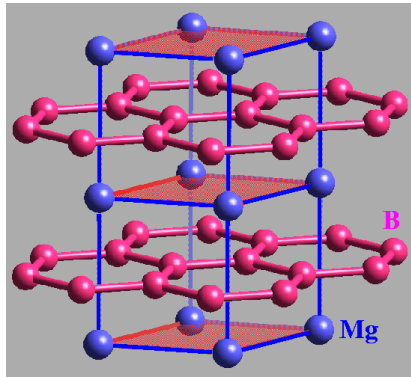
	Single-gap	multi-gap
$\lambda$	$\lambda \sim 0.6$	$\lambda_{max} \sim 1$
$\omega_{ph}$	75 meV	75 meV
$T_c$	19 K	39 K

**Small Impurity interband scattering**



**Cooper pairs are not mixed maintaining their characteristics with two perfectly separated gaps**





$\sigma$  states are 2D, confined in the B planes, strongly coupled with  $E_{2g}$   
 $\pi$  states are 3D extending in all directions, weakly coupled with phonons

Different types of impurities can affect selectively  $\sigma$  and  $\pi$  mean free paths with distinct consequences on the transport properties:  
 the most deeply investigated are **Al on Mg** and **C on B**.

Both substitutions **reduce  $T_c$** : **e-doping** reduces the DOS ( $\sigma$  band filling)

**Al on Mg** affects only the  $\pi$  mean free path:

$\pi$  bands are in “dirty limit”

$\sigma$  bands are in “clean limit”

**C on B** affects the  $\pi$  and  $\sigma$  mean free path:

$\pi$  and  $\sigma$  bands are in “dirty limit”

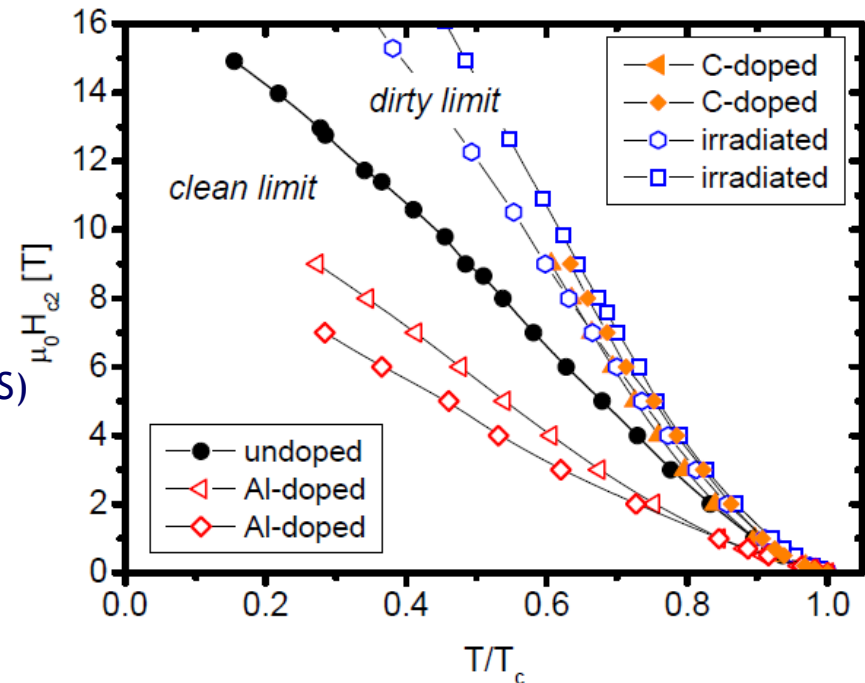
**Al doped samples** are in clean limit and  $H_{c2}$  **decreases**  
 (weakening of the coupling due to the decrease of the DOS)

**C doped samples,**

as well as **irradiated samples,** are in dirty limit

and  $H_{c2}$  **increases**

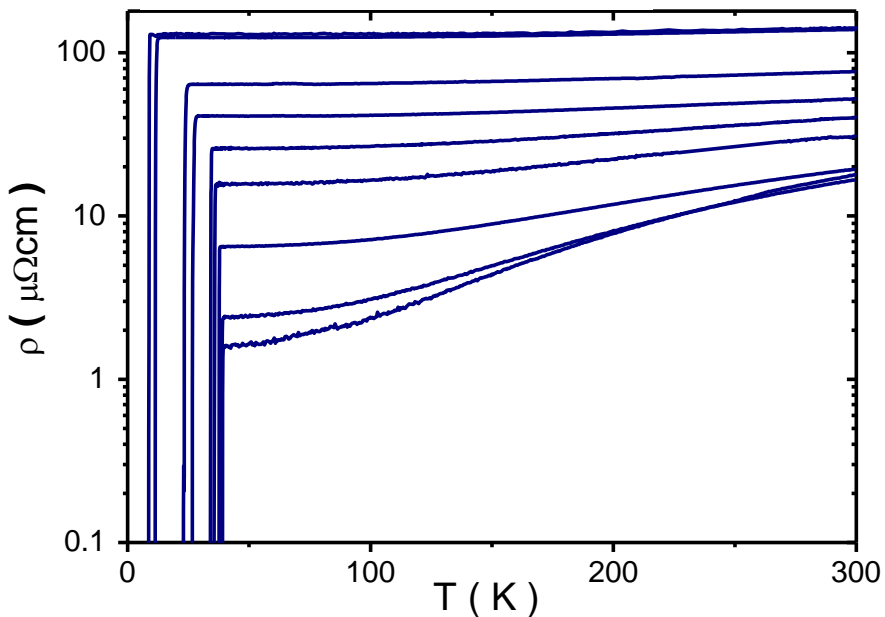
(decrease of the DOS but with a decrease of the mfp)





**Neutron irradiation:** a way to introduce defect without doping demonstrating the role of interband scattering in **suppress  $T_c$**  and **producing the merging of the two energy gaps**

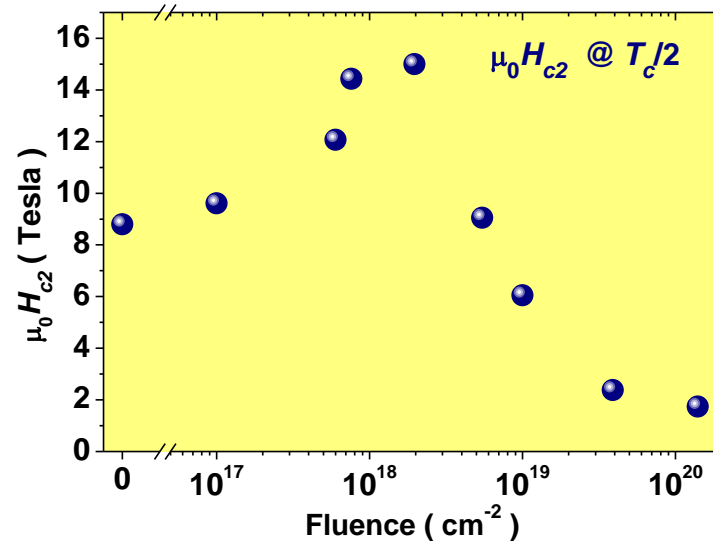
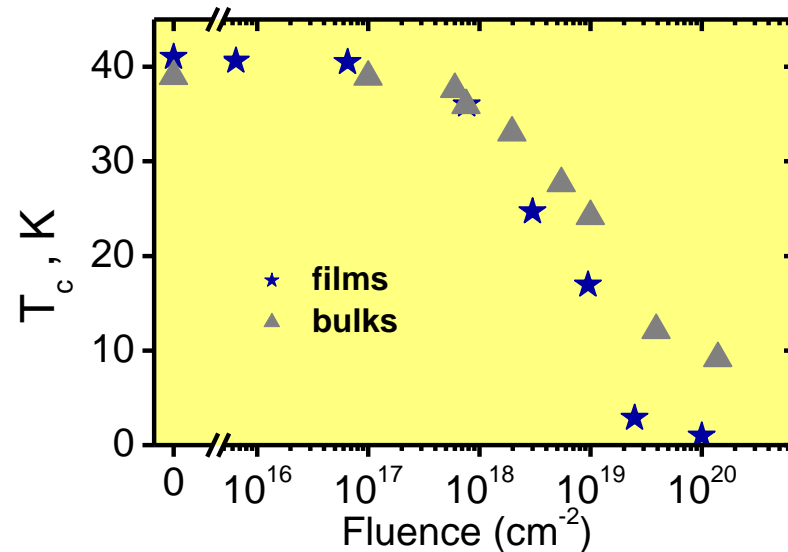
Tarantini et al.  
PHYSICAL REVIEW B 73, 134518 (2006)

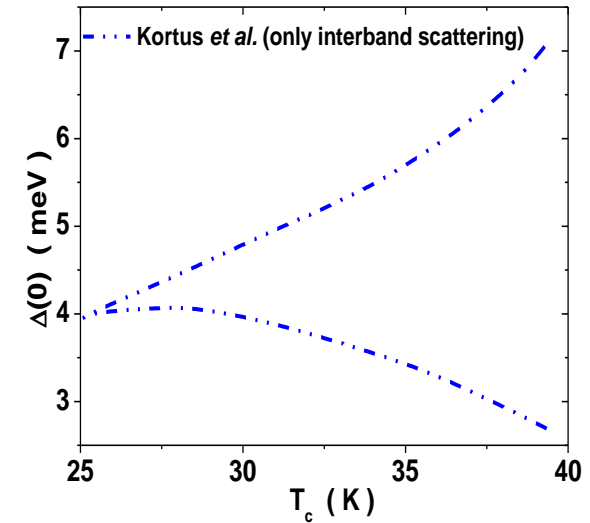
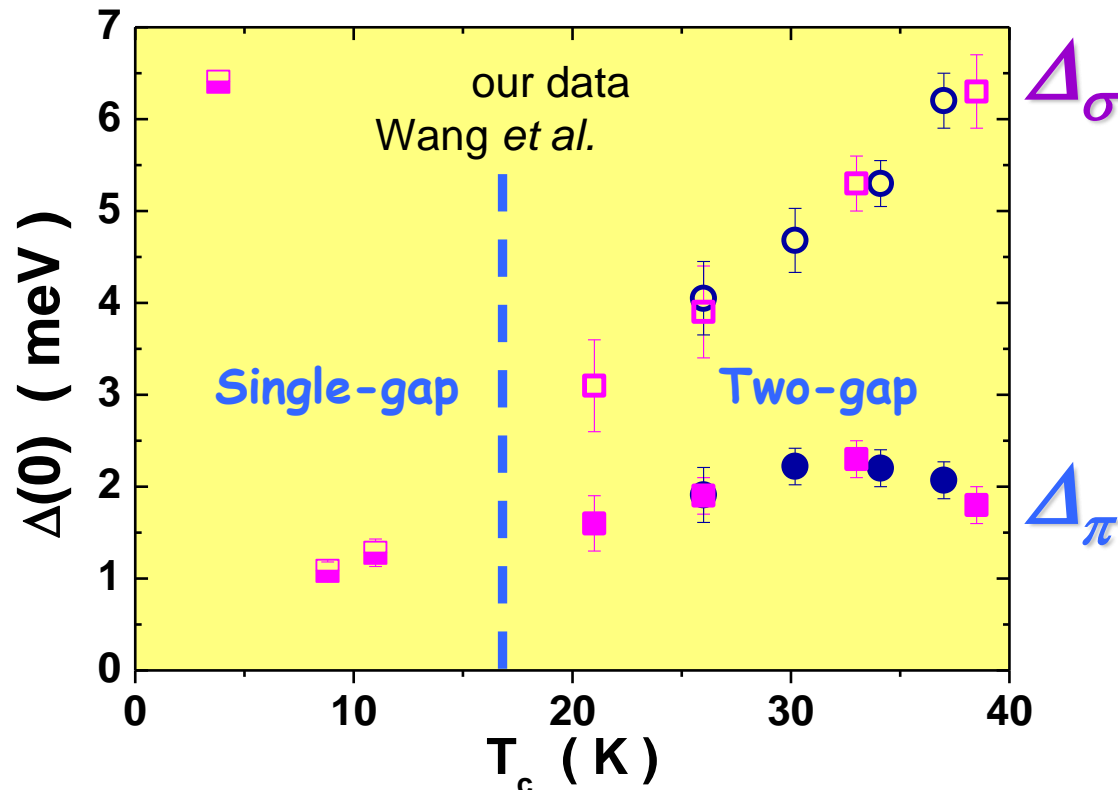


$T_c$  decreases down to 6K, resistivity increases

$H_{c2}$  strongly increases up to fluence of  $2 \times 10^{18} \text{ cm}^{-2}$ , while  $T_c$  is only slightly decreased.

At higher fluence, both  $T_c$  and  $H_{c2}$  are strongly suppressed





The cross-over from two- to single-gap behaviour occurs between 21 and 11 K

$\Delta_\sigma$  decreases linearly with irradiation

Decreasing  $T_c$ ,  $\Delta_\pi$  increases down to  $\sim 30$  K, then smoothly decreases



**Interband scattering**

M. Putti, C. Tarantini *et al.* *Physical Review Letters* **96**, 077003 (2006)

D. Daghero, C. Tarantini *et al.*, *Physical Review B* **74**, 174519 (2006)

## TOPICAL REVIEW

# The widely variable resistivity of MgB<sub>2</sub> samples

John M Rowell<sup>1</sup>

- T<sub>c</sub> ~ 39K in samples with very low resistivity (10 μΩ cm at 300K and 0.5 μΩ cm at 40K)
- T<sub>c</sub> ~ 39K in samples with high resistivity
- The behavior of rho Vs T is always T<sup>2</sup> to T<sup>3</sup>

No correlation between rho and T<sub>c</sub>?

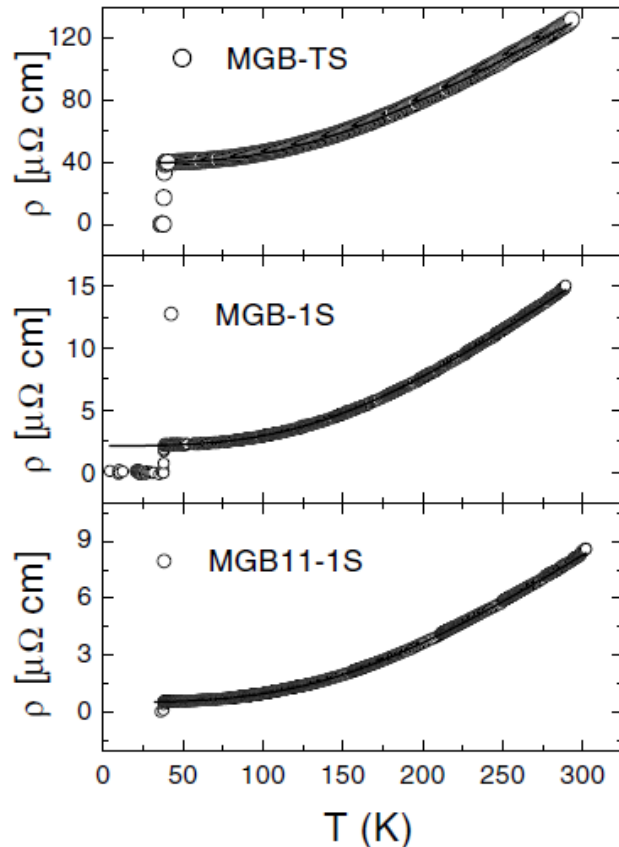
Mazin:

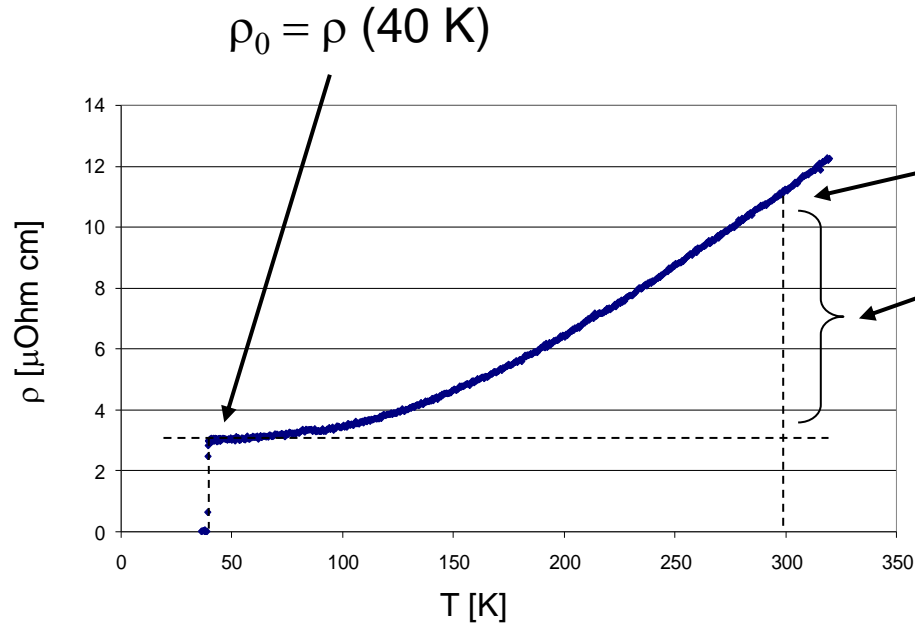
Evidence of band disparity,

No interband scattering, No suppression of T<sub>c</sub>

Rowell:

the resistivity increase in most cases could be simply ascribed to a **weak connection between the grains** which causes a reduction in the **effective current-carrying cross-sectional area** of the sample.





$\rho$  (300 K)

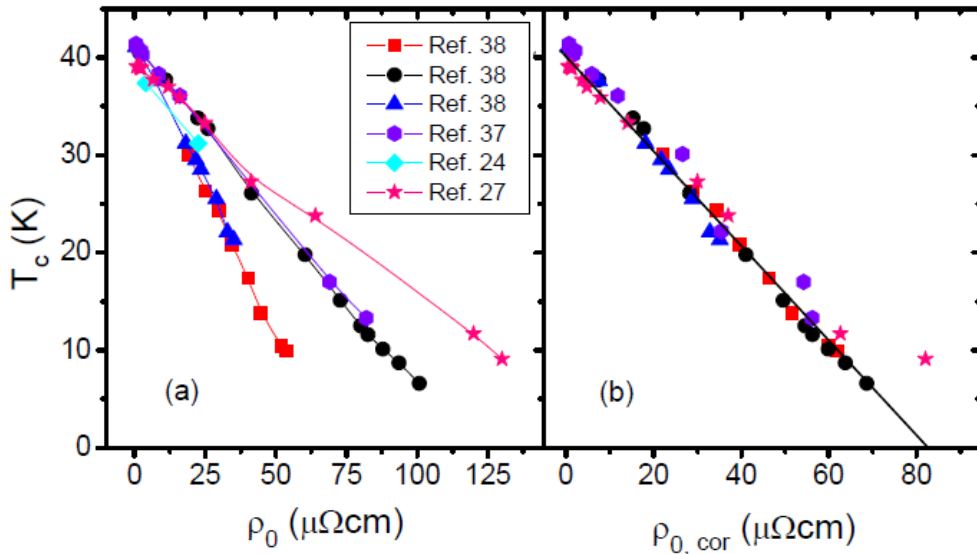
$\Delta\rho = \rho(300\text{ K}) - \rho(40\text{ K})$

Assuming  $\Delta\rho = \text{const}$ ,  $\Delta\rho_g = 7.5\ \mu\text{Ohm cm}$   
like in the case of perfect connectivity

$$\rho_{0,g} = \frac{\Delta\rho_g}{\Delta\rho} \rho_0$$

## Effective area fraction

$$A_f = \frac{\Delta\rho_g}{\Delta\rho}$$



$T_c$  versus residual resistivity of several series of irradiated films or bulk:  
All data collapse on the same line using the “Rowell” normalized resistivity