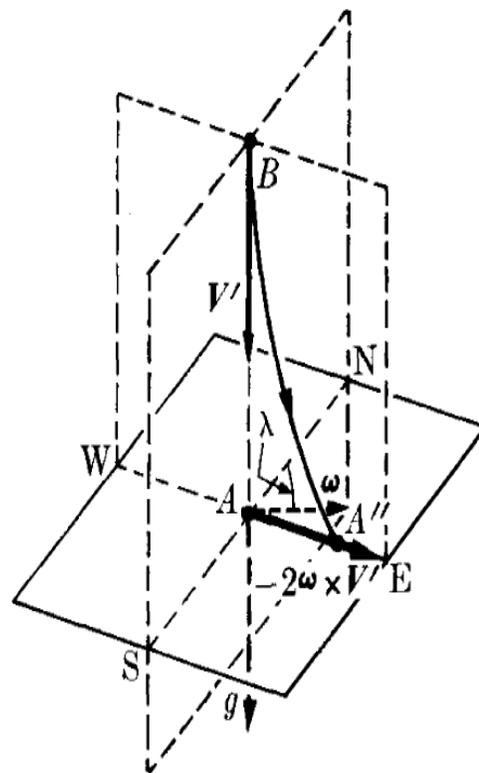


(a) Northern hemisphere



(b) Southern hemisphere

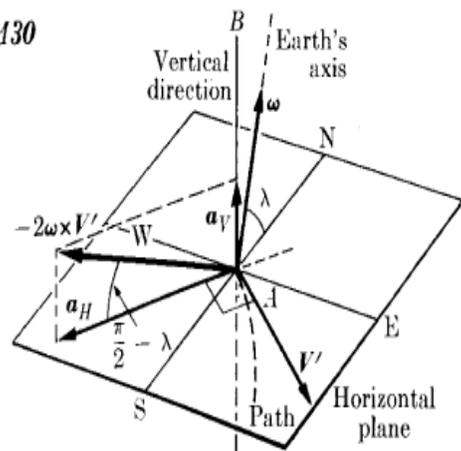
**Fig. 6-10.** Deviation to the east in the Northern (Southern) hemisphere of a falling body due to Coriolis acceleration.

Il moto di un corpo rispetto alla Terra:  $V' \neq 0$

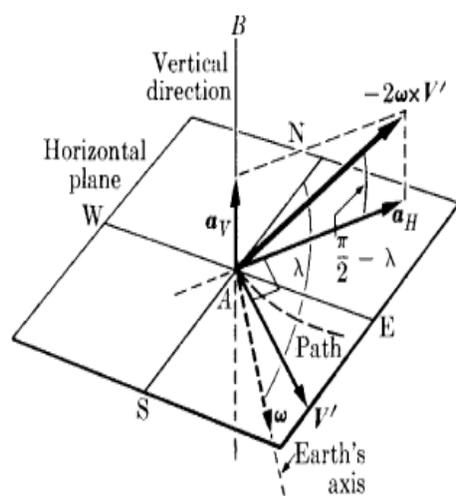
$$\mathbf{g} = \mathbf{g}_0 - \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) - 2\boldsymbol{\omega} \times \mathbf{V}'$$

Il termine di Coriolis introduce una deviazione verso est nel moto in caduta libera

Combining this Coriolis effect with the centrifugal effect, the body will fall on a point southeast of  $A$  in the Northern hemisphere and northeast of  $A$  in the Southern hemisphere. This effect, which is negligible in most cases, must be carefully taken into account both in high-altitude bombing and in intercontinental ballistic missiles. Coriolis acceleration also seriously affects the paths of rockets and of satellites, due to their great velocities.

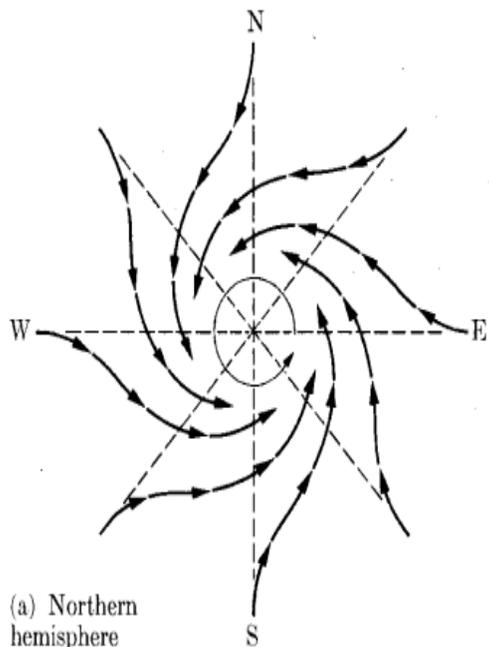


(a) Northern hemisphere

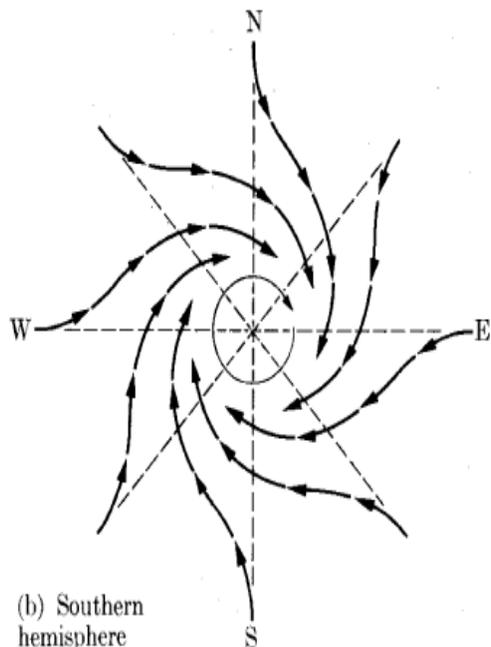


(b) Southern hemisphere

**Fig. 6-11.** Coriolis acceleration. When a body moves in a horizontal plane, the horizontal component of the Coriolis acceleration points to the right (left) of the direction of motion in the Northern (Southern) hemisphere. Here  $V'$  is in the horizontal plane,  $\omega$  is in the plane defined by  $AB$  and  $NS$ , and  $a_H$  is perpendicular to  $V'$ .



(a) Northern hemisphere



(b) Southern hemisphere

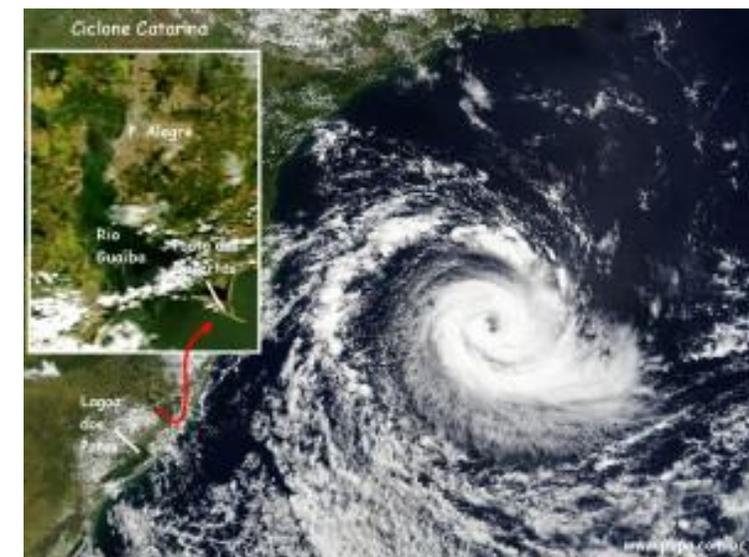
Il moto di un corpo rispetto alla Terra:  $V' \neq 0$   
gli uragani

$$\mathbf{g} = \mathbf{g}_0 - \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) - 2\boldsymbol{\omega} \times \mathbf{V}'$$

Katrina



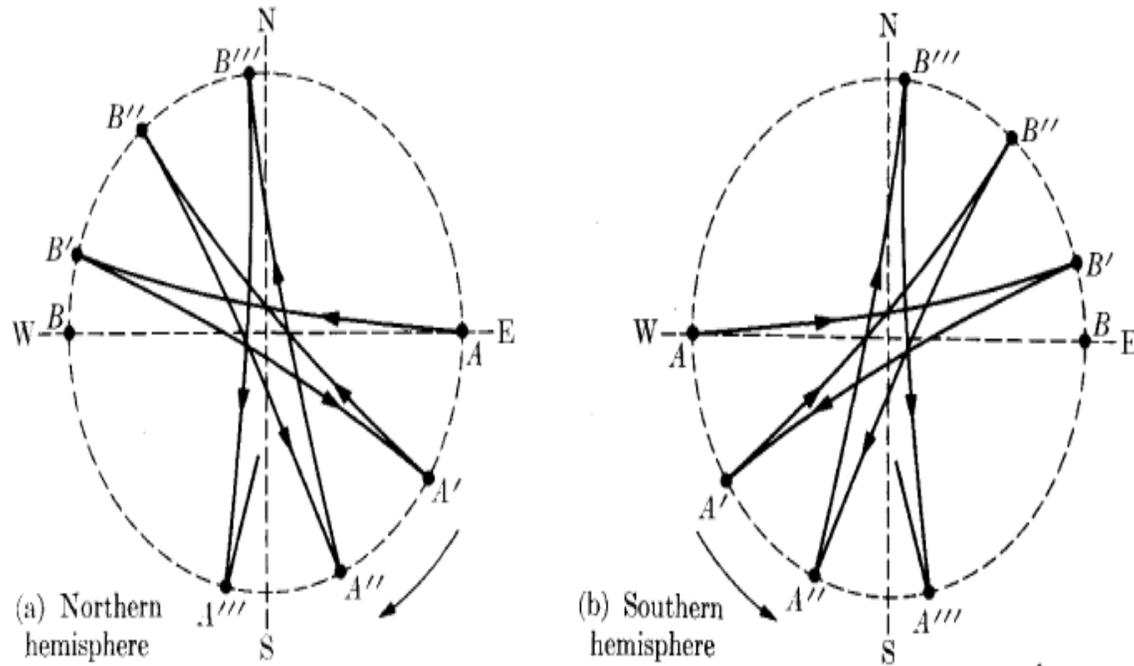
Catarina



# Il moto di un corpo rispetto alla Terra: $V' \neq 0$

## Il pendolo

$$\mathbf{g} = \mathbf{g}_0 - \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) - 2\boldsymbol{\omega} \times \mathbf{V}'$$



**Fig. 6-13.** Rotation of plane of oscillation of pendulum as a result of Coriolis acceleration. (Rotation in the Southern hemisphere is in the opposite direction to that in Northern.)

As a second example, let us consider the oscillations of a pendulum. When the amplitude of the oscillations is small, we can assume that the motion of the bob is along a horizontal path. If the pendulum were initially set to oscillate in the east-west direction and were released at  $A$  (see Fig. 6-13), it would continue oscillating between  $A$  and  $B$  if the earth were not rotating. But because of the Coriolis acceleration due to the earth's rotation, the path of the pendulum is deflected continuously to the right in the Northern hemisphere and to the left in the Southern hemisphere. Therefore, at the end of the first oscillation, it reaches  $B'$  instead of  $B$ . On its return, it goes to  $A'$  and not to  $A$ . Therefore, in successive complete oscillations, it arrives at  $A''$ ,  $A'''$ , etc. In other words, the plane of oscillation of the pendulum rotates clockwise in the Northern hemisphere and counterclockwise in the Southern hemisphere. We leave to the student the verification of the fact that the angle through which the plane of oscillation rotates each hour is  $15^\circ \sin \lambda$ . The effect has been much exaggerated in Fig. 6-13; it is maximum at the poles and zero at the equator.

# Rotazione del piano di oscillazione del pendolo, calcolo dell'angolo di rotazione.

Il pendolo parte da fermo spostato dalla posizione di equilibrio (punto E). Ogni mezza periodo il piano del pendolo ruota di una stessa quantità. La forza che agisce sul pendolo in direzione S-N è:

$$F = ma = 2 m \omega V(t) \sin \lambda - m \omega_p^2 y$$

Cioè la somma di forza di Coriolis e forza peso che tende a richiamarlo verso la linea E-O.  $V(t)$  è la velocità del pendolo,  $y$  è la coordinate lungo la direzione S-N. Si ha:

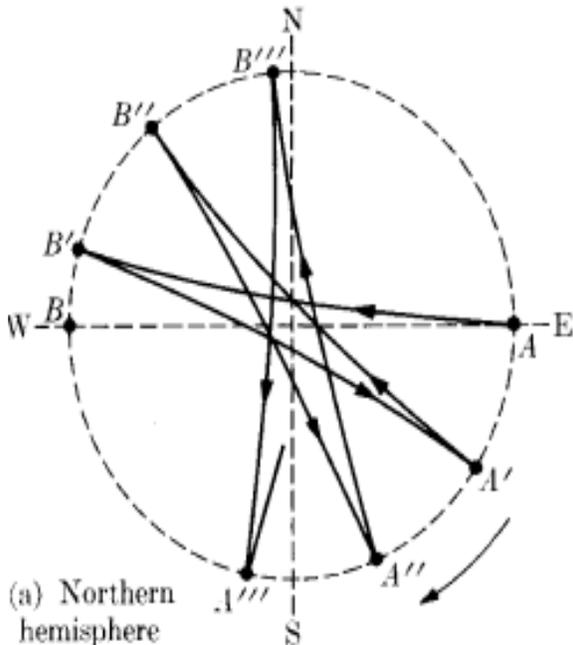
$$V(t) = V_{max} \sin(\omega_p t)$$

con  $\omega_p$  velocità angolare del pendolo e  $V_{max}$  sua velocità massima. Poiché le correzioni non inerziali sono piccole rispetto alla forza di gravità,  $\omega_p$  e  $V_{max}$  sono quelle del pendolo inerziale.

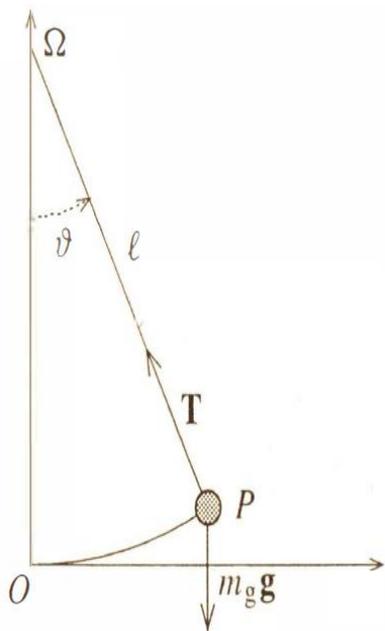
$$\omega_p = \sqrt{\frac{g}{l}}; \quad V_{max} = \theta_{max} \sqrt{g l}$$

$\theta_{max}$  è l'ampiezza dell'angolo in E. Dobbiamo risolvere la seguente equazione differenziale con le condizioni iniziali per la coordinata  $y$  riportate a fianco.

$$\frac{d^2 y}{dt^2} = 2 \omega V_{max} \sin(\omega_p t) \sin \lambda - \omega_p^2 y; \quad y(0) = 0; \quad y'(0) = 0$$



(a) Northern hemisphere



# Rotazione del piano di oscillazione del pendolo, calcolo dell'angolo di rotazione.

Si ottiene

$$y(t) = \frac{\omega V_{max} \sin \lambda}{\omega_p^2} (-\omega_p t \cos(\omega_p t) + \sin(\omega_p t))$$

Una semi oscillazione del pendolo (distanza tra E e W) impiega un tempo paria a metà del periodo del pendolo, in tale tempo lo spostamento a destra è dato da

$$y\left(\frac{\pi}{\omega_p}\right) = \frac{\pi \omega V_{max} \sin \lambda}{\omega_p^2} = \frac{\pi}{\omega_p} \theta_{max} l \omega \sin \lambda = \frac{T_p}{2} \theta_{max} l \omega \sin \lambda$$

L'angolo di rotazione del piano del pendolo in un semi periodo  $\frac{T_p}{2}$  è dato da

$$\alpha = \frac{s}{l \theta_{max}} = \omega \sin \lambda \frac{T_p}{2}$$

Essendoci tantissimi semi periodi di oscillazione in un'ora, in pratica, l'angolo di rotazione del piano del pendolo è proporzionale al tempo in cui il pendolo oscilla. Se oscilla per un'ora

$$\alpha = \omega \sin \lambda \ 1 \text{ ora} = \frac{2\pi}{86.400 \text{ s}} \sin \lambda \ 3600 \text{ s} = \frac{2\pi}{24} \sin \lambda \text{ rad} = 15^\circ \sin \lambda$$

